

HARMONIC ANALYSIS OF ADDITIVE LÉVY PROCESSES

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ABSTRACT. Let X_1, \dots, X_N denote N independent d -dimensional Lévy processes, and consider the N -parameter random field

$$\mathfrak{X}(t) := X_1(t_1) + \dots + X_N(t_N).$$

First we demonstrate that for all nonrandom Borel sets $F \subseteq \mathbf{R}^d$, the Minkowski sum $\mathfrak{X}(\mathbf{R}_+^N) \oplus F$, of the range $\mathfrak{X}(\mathbf{R}_+^N)$ of \mathfrak{X} with F , can have positive d -dimensional Lebesgue measure if and only if a certain capacity of F is positive. This improves our earlier joint effort with Yuquan Zhong (2003) by removing a symmetry-type condition there. Moreover, we show that under mild regularity conditions, our necessary and sufficient condition can be recast in terms of one-potential densities. This rests on developing results in classical [non-probabilistic] harmonic analysis that might be of independent interest. As was shown in Khoshnevisan, Xiao, and Zhong (2003), the potential theory of the type studied here has a large number of consequences in the theory of Lévy processes. We present a few new consequences here.

1. Introduction

1.1. Background. It is known that, for all integers $d \geq 2$, the range of d -dimensional Brownian motion has zero Lebesgue measure. See Lévy (1940) for $d = 2$, Ville (1942) for $d = 3$, and Kakutani (1944a) for the remaining assertions. There is a perhaps better-known, but equivalent, formulation of this theorem: When $d \geq 2$, the range of d -dimensional Brownian motion does not hit points. Kakutani (1944b) has generalized this by proving that, for all integers $d \geq 1$, the range of d -dimensional Brownian motion can hit a nonrandom Borel set $F \subseteq \mathbf{R}^d$ if and only if $\text{cap}(F) > 0$, where cap denotes, temporarily, the logarithmic capacity if $d = 2$ and the Riesz capacity of index $d - 2$ if $d \geq 3$; the case $d = 1$ is elementary. [Actually, Kakutani's paper discusses only the planar case. The theorem, for $d \geq 3$, is pointed out in Dvoretzky, Erdős, and Kakutani (1950).]

Kakutani's theorem is the starting point of a deep *probabilistic potential theory* initiated by Hunt (1958, 1957a, 1957b, 1956). The literature on this topic is rich and quite large; see, for

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example, the books by Blumenthal and Gettoor (1968), Doob (2001), Fukushima, Ōshima, and Takeda (1994), Gettoor (1990), and Röckner (1993), together with their combined bibliography.

One of the central assertions of probabilistic potential theory is that a nice Markov process will hit a nonrandom measurable set F if and only if $\text{cap}(F) > 0$, where cap is a certain natural capacity in the sense of G. Choquet (Dellacherie and Meyer, 1978, Chapter III, pp. 51–55). Moreover, that capacity is defined solely, and fairly explicitly, in terms of the Markov process itself.

There are interesting examples where F is itself random. For instance, suppose X is d -dimensional standard Brownian motion, and $F = Y((0, \infty))$ is the range—minus the starting point—of an independent standard Brownian motion Y on \mathbf{R}^d . In this particular case, it is well known that

$$(1.1) \quad \mathbf{P} \{X(s) = Y(t) \text{ for some } s, t > 0\} > 0 \quad \text{if and only if} \quad d \leq 3.$$

This result was proved by Lévy (1940) for $d = 2$, Kakutani (1944a) for $d \geq 5$, and Dvoretzky, Erdős, and Kakutani (1944b) for $d = 3, 4$. Peres (1996b, 1996a) and Khoshnevisan (2003) contain different elementary proofs of this fact.

There are many generalizations of (1.1) in the literature. For example, Dvoretzky, Erdős, and Kakutani (1944b) proved that the paths of an arbitrary number of independent planar Brownian motions can intersect. While LeGall (1987) proved that the trajectories of a planar Brownian motion can intersect itself countably many times. And Dvoretzky, Erdős, Kakutani, and Taylor (1957) showed that three independent Brownian-motion trajectories in \mathbf{R}^d can intersect if and only if $d \leq 2$. For other results along these lines, see Hawkes (1978b, 1977, 1976/77), Hendricks (1979, 1973/74), Kahane (1985, Chapter 16, Section 6), Lawler (1989, 1985, 1982), Pemantle, Peres, and Shapiro (1996), Le Gall (1992), Peres (1999, 1996b, 1996a), Rogers (1989), Tongring (1988), and their combined bibliography.

For a long time, a good deal of effort was concentrated on generalizing (1.1) to other concrete Markov processes than Brownian motion. But the problem of deciding when the paths of N independent (general but nice) Markov processes can intersect remained elusive. It was finally settled, more than three decades later, by Fitzsimmons and Salisbury (1989), whose approach was to consider the said problem as one about a certain multiparameter Markov process.

To be concrete, let us consider the case $N = 2$, and let $X := \{X(t)\}_{t \geq 0}$ and $Y := \{Y(t)\}_{t \geq 0}$ denote two independent (nice) Markov processes on a (nice) state space S . The starting point of the work of Fitzsimmons and Salisbury is the observation that $\mathbf{P}\{X(s) = Y(t) \text{ for some } s, t > 0\} > 0$ if and only if the two-parameter Markov process $X \otimes Y$ hits the

diagonal $\text{diag } S := \{x \otimes x : x \in S\}$ of $S \times S$, where

$$(1.2) \quad (X \otimes Y)(s, t) := \begin{pmatrix} X(s) \\ Y(t) \end{pmatrix} \quad \text{for all } s, t \geq 0.$$

In the special case that X and Y are Lévy processes, the Fitzsimmons–Salisbury theory was used to solve the then-long-standing Hendricks–Taylor conjecture (Hendricks and Taylor, 1979).

The said connection to multiparameter processes is of paramount importance in the Fitzsimmons–Salisbury theory, and appears earlier in the works of Evans (1987a, 1987b). See also Le Gall, Rosen, and Shieh (1989) and Salisbury (1996, 1992, 1988). Walsh (1986, pp. 364–368) discusses a connection between $X \otimes Y$ and the Dirichlet problem for the biLaplacian $\Delta \otimes \Delta$ on the bi-disc of $\mathbf{R}^d \times \mathbf{R}^d$.

The Fitzsimmons–Salisbury theory was refined and generalized in different directions by Hirsch (1995), Hirsch and Song (1999, 1996, 1995a, 1995b, 1995c, 1995d, 1994), and Khoshnevisan (2002, 1999). See also Ren (1990) who derives an implicit-function theorem in classical Wiener space by studying a very closely-related problem.

The two-parameter process $X \otimes Y$ itself was introduced earlier in the works of Wolpert (1978), who used $X \otimes Y$ to build a $(\phi^\kappa)_2$ model of the Euclidean field theory. This too initiated a very large body of works. For some of the earlier examples, see the works by Aizenman (1985), Albeverio and Zhou (1996), Dynkin (1987, 1986, 1985, 1984a, 1984b, 1984c, 1983a, 1983b, 1981, 1980), Felder and Fröhlich (1985), Rosen (1983, 1984), and Westwater (1980, 1981, 1982). [This is by no means an exhaustive list.]

In the case that X and Y are Lévy processes on \mathbf{R}^d [i.e., have stationary independent increments], $X \otimes Y$ is an example of the so-called additive Lévy processes. But as it turns out, it is important to maintain a broader perspective and consider more than two Lévy processes. With this in mind, let X_1, \dots, X_N denote N independent Lévy processes on \mathbf{R}^d such that each X_j is normalized via the Lévy–Khintchine formula (Bertoin, 1996; Sato, 1999):

$$(1.3) \quad \mathbb{E} \exp(i\xi \cdot X_j(t)) = \exp(-t\Psi_j(\xi)) \quad \text{for all } t \geq 0 \text{ and } \xi \in \mathbf{R}^d.$$

The function Ψ_j is called the *characteristic exponent*—or *Lévy exponent*—of X_j , and its defining property is that Ψ_j is a negative-definite function (Schoenberg, 1938; Berg and Forst, 1975).

1.2. The main results. The main object of this paper is to develop some basic probabilistic potential theory for the following N -parameter random field \mathfrak{X} with values in \mathbf{R}^d :

$$(1.4) \quad \mathfrak{X}(\mathbf{t}) := X_1(t_1) + \dots + X_N(t_N) \quad \text{for all } \mathbf{t} := (t_1, \dots, t_N) \in \mathbf{R}_+^N.$$

On a few occasions we might write $(\oplus_{j=1}^N X_j)(\mathbf{t})$ in place of $\mathfrak{X}(\mathbf{t})$, as well.

The random field \mathfrak{X} is a so-called *additive Lévy process*, and is characterized by its multi-parameter Lévy–Khintchine formula:

$$(1.5) \quad \mathbb{E} \exp(i\xi \cdot \mathfrak{X}(\mathbf{t})) = \exp(-\mathbf{t} \cdot \Psi(\xi)) \quad \text{for all } \mathbf{t} \in \mathbf{R}_+^N \text{ and } \xi \in \mathbf{R}^d;$$

where $\Psi(\xi) := (\Psi_1(\xi), \dots, \Psi_N(\xi))$ is the *characteristic exponent* of \mathfrak{X} . Our goal is to describe the potential-theoretic properties of \mathfrak{X} solely in terms of its characteristic exponent Ψ . Thus, it is likely that our harmonic-analytic viewpoint can be extended to study the potential theory of more general multiparameter Markov processes that are based on the Feller processes of Jacob (2005, 2002, 2001).

In order to describe our main results let us first consider the kernel

$$(1.6) \quad K_\Psi(\xi) := \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi)} \right) \quad \text{for all } \xi \in \mathbf{R}^d.$$

When $N = 1$, this kernel plays a central role in the works of Orey (1967) and Kesten (1969).

The kernel for general N was introduced first by Evans (1987b); see also Khoshnevisan, Xiao, and Zhong (2005).

Based on the kernel K_Ψ , we define, for all Schwartz distributions μ on \mathbf{R}^d ,

$$(1.7) \quad I_\Psi(\mu) := \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} |\hat{\mu}(\xi)|^2 K_\Psi(\xi) d\xi.$$

We are primarily interested in the case where μ is a real-valued locally integrable function, or a σ -finite Borel measure on \mathbf{R}^d . In either case, we refer to $I_\Psi(\mu)$ as the *energy* of μ . Our notion of energy corresponds to a *capacity* cap_Ψ , which is the following set function: For all Borel sets $F \subseteq \mathbf{R}^d$,

$$(1.8) \quad \operatorname{cap}_\Psi(F) := \left[\inf_{\mu \in \mathcal{P}_c(F)} I_\Psi(\mu) \right]^{-1},$$

where $\mathcal{P}_c(F)$ denotes the collection of all compactly-supported Borel probability measures on F , $\inf \emptyset := \infty$, and $1/\infty := 0$.

The following is the first central result of this paper. Here and throughout, λ_k denotes k -dimensional Lebesgue measure on \mathbf{R}^k for all integers $k \geq 1$.

Theorem 1.1. *Let \mathfrak{X} be an N -parameter additive Lévy process on \mathbf{R}^d with exponent Ψ . Then, for all Borel sets $F \subseteq \mathbf{R}^d$,*

$$(1.9) \quad \mathbb{E} [\lambda_d(\mathfrak{X}(\mathbf{R}_+^N) \oplus F)] > 0 \quad \text{if and only if} \quad \operatorname{cap}_\Psi(F) > 0.$$

Remark 1.2. (1) Theorem 1.1 in the one-parameter setting is still very interesting, but much easier to derive. See Kesten (1969) for the case that $F := \{0\}$ and Hawkes (1984) for general F . For a scholarly pedagogic account see Bertoin (1996, p. 60).

- (2) One can view Theorem 1.1 as a contribution to the theory of Dirichlet forms for a class of infinite-dimensional Lévy processes. These Lévy processes are in general non-symmetric. Röckner (1993) describes a general theory of Dirichlet forms for nice infinite-dimensional Markov processes that are not necessarily symmetric. It would be interesting to know if the processes of the present paper lend themselves to the analysis of the general theory of Dirichlet forms. We have no conjectures along these lines. \square

Our earlier collaborative effort with Yuquan Zhong (2003) yielded the conclusion of Theorem 1.1 under an exogenous technical condition on X_1, \dots, X_N . A first aim of this paper is to establish the fact that Theorem 1.1 holds in complete generality. Also, we showed in our earlier works (Khoshnevisan and Xiao, 2005; , 2004, Khoshnevisan, Xiao, and Zhong ; 2003) that such a theorem has a large number of consequences, many of them in the classical theory of Lévy processes itself. Next we describe a few such consequences that are nontrivial due to their intimate connections to harmonic analysis.

Our next result provides a criterion for a Borel set $F \subseteq \mathbf{R}^d$ to contain intersection points of N independent Lévy processes. It completes and complements the well-known results of Fitzsimmons and Salisbury (1989). See also Corollary 9.3 and Remark 9.4 below.

Theorem 1.3. *Let X_1, \dots, X_N be independent Lévy processes on \mathbf{R}^d , and assume that each X_j has a one-potential density $u_j : \mathbf{R}^d \rightarrow \bar{\mathbf{R}}_+$ such that $u_j(0) > 0$. Then, for all nonempty Borel sets $F \subseteq \mathbf{R}^d$,*

$$(1.10) \quad \mathbf{P} \{X_1(t_1) = \dots = X_N(t_N) \in F \text{ for some } t_1, \dots, t_N > 0\} > 0$$

if and only if there exists a compact-support Borel probability measure μ on F such that

$$(1.11) \quad \int_{\mathbf{R}^d} \dots \int_{\mathbf{R}^d} |\hat{\mu}(\xi^1 + \dots + \xi^N)|^2 \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi^j)} \right) d\xi^1 \dots d\xi^N < \infty.$$

Suppose, in addition, that every u_j is continuous on \mathbf{R}^d , and finite on $\mathbf{R}^d \setminus \{0\}$. Then, another equivalent condition is that there exists a compact-support probability measure μ on F such that

$$(1.12) \quad \iint \prod_{j=1}^N \left(\frac{u_j(x-y) + u_j(y-x)}{2} \right) \mu(dx) \mu(dy) < \infty.$$

In order to describe our next contribution, let us recall that the *one-potential measure* U of a Lévy process $X := \{X(t)\}_{t \geq 0}$ on \mathbf{R}^d is defined as

$$(1.13) \quad U(A) := \int_0^\infty \mathbf{P} \{X(t) \in A\} e^{-t} dt,$$

for all Borel sets $A \subseteq \mathbf{R}^d$. Next we offer a two-parameter “additive variant” which requires fewer technical conditions than Theorem 1.3.

Theorem 1.4. *Suppose X_1 and X_2 are independent Lévy processes on \mathbf{R}^d with respective one-potential measures U_1 and U_2 . Suppose $U_1(dx) = u_1(x) dx$, where $u_1 : \mathbf{R}^d \rightarrow \bar{\mathbf{R}}_+$, and $u_1 * U_2 > 0$ almost everywhere. Then, for all Borel sets $F \subseteq \mathbf{R}^d$,*

$$(1.14) \quad \mathbf{P} \{X_1(t_1) + X_2(t_2) \in F \text{ for some } t_1, t_2 > 0\} > 0$$

if and only if there exists a compact-support Borel probability measure μ on F such that

$$(1.15) \quad \int_{\mathbf{R}^d} |\hat{\mu}(\xi)|^2 \operatorname{Re} \left(\frac{1}{1 + \Psi_1(\xi)} \right) \operatorname{Re} \left(\frac{1}{1 + \Psi_2(\xi)} \right) d\xi < \infty.$$

Suppose, in addition, that u_1 is continuous on \mathbf{R}^d , and finite on $\mathbf{R}^d \setminus \{0\}$. Then, (1.14) holds if and only if there exists a compact-support probability measure μ on F such that

$$(1.16) \quad \iint Q(x - y) \mu(dx) \mu(dy) < \infty,$$

where

$$(1.17) \quad Q(x) := \int_{\mathbf{R}^d} \left[\frac{u_1(x + y) + u_1(x - y) + u_1(-x + y) + u_1(-x - y)}{4} \right] U_2(dy)$$

for all $x \in \mathbf{R}^d$.

Among other things, Theorem 1.4 confirms a conjecture of Bertoin (1999b) and Bertoin (1999a, p. 49); see Remark 8.1 for details.

Finally we mention a result on the Hausdorff dimension of the set of intersections of the sample paths of Lévy processes.

Theorem 1.5. *Let X_1, \dots, X_N be independent Lévy processes on \mathbf{R}^d , and assume that each X_j has a one-potential density $u_j : \mathbf{R}^d \rightarrow \bar{\mathbf{R}}_+$ such that $u_j(0) > 0$. Then, almost surely on $\{\cap_{k=1}^N X_k(\mathbf{R}_+) \neq \emptyset\}$,*

$$(1.18) \quad \begin{aligned} & \dim_{\mathbf{H}} \bigcap_{k=1}^N X_k(\mathbf{R}_+) \\ &= \sup \left\{ s \in (0, d) : \int_{(\mathbf{R}^d)^N} \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi^j)} \right) \frac{d\xi}{1 + \|\xi^1 + \dots + \xi^N\|^{d-s}} < \infty \right\}, \end{aligned}$$

where $\sup \emptyset := 0$. Suppose, in addition, that the u_j 's are continuous on \mathbf{R}^d , and finite on $\mathbf{R}^d \setminus \{0\}$. Then, almost surely on $\{\cap_{k=1}^N X_k(\mathbf{R}_+) \neq \emptyset\}$,

$$(1.19) \quad \dim_{\mathbf{H}} \bigcap_{k=1}^N X_k(\mathbf{R}_+) = \sup \left\{ s \in (0, d) : \int_{\mathbf{R}^d} \prod_{j=1}^N \left(\frac{u_j(z) + u_j(-z)}{2} \right) \frac{dz}{\|z\|^s} < \infty \right\}.$$

In the remainder of the paper we prove Theorem 1.1 and its stated corollaries in the order in which they are presented. Finally, we conclude by two zero-one laws for the Lebesgue measure and capacity of the range of an additive Lévy process, that, we believe, might have independent interest.

We end this section with four problems and conjectures.

Open problem 1. Throughout this paper, we impose continuity conditions on various one-potential densities. This is mainly because we are able to develop general harmonic-analytic results only for kernels that satisfy some regularity properties. Can the continuity conditions be dropped? We believe the answer is “yes.” This is motivated, in part, by the following fact, which follows from inspecting the proofs: The condition “ u is continuous on \mathbf{R}^d and finite on $\mathbf{R}^d \setminus \{0\}$ ” is used only for proving the “if” portions in the second parts of Theorems 1.3 and 1.4.

Open problem 2. Jacob (2005; 2002; 2001) has constructed a very large class of Feller processes that behave locally like Lévy processes. Moreover, his construction is deeply connected to harmonic analysis. Because the results of the present paper involve mainly the local structure of Lévy processes, and are inextricably harmonic analytic, we ask: Is it possible to study the harmonic-analytic potential theory of several Jacob processes by somehow extending the methods of the present paper?

Open problem 3. We ask: Is there a “useful” theory of excessive functions and/or measures for additive Lévy processes (or more general multiparameter Markov processes)? This question is intimately connected to Open Problem 1, but deserves to be asked on its own. In the one-parameter case, the answer is a decisive “yes” (Gettoor, 1990). But the one-dimensional theory does not appear to readily have a suitable extension to the multiparameter setting.

Open problem 4. We conjecture that, under the conditions of Theorem 1.5, the following holds almost surely on $\{\cap_{k=1}^N X_k(\mathbf{R}_+) \neq \emptyset\}$:

$$(1.20) \quad \dim_{\mathbf{H}} \bigcap_{k=1}^N X_k(\mathbf{R}_+) = \sup \left\{ s \in (0, d) : \int_{(-1,1)^d} \prod_{j=1}^N \left(\frac{u_j(z) + u_j(-z)}{2} \right) \frac{dz}{\|z\|^s} < \infty \right\}.$$

[The difference between this and (1.19) is in the range of the integrals.] But in all but one case we have no proof; see Remark 9.6 below for the mentioned case. As we shall see in that remark, what we actually prove is the following harmonic-analytic fact: *Suppose u is the one-potential density of a Lévy process, $u(0) > 0$, u is continuous on \mathbf{R}^d , and u is finite on $\mathbf{R}^d \setminus \{0\}$. Then the local square-integrability of u implies the [global] square-integrability of u .* We believe that the following more general result holds: If u_1, \dots, u_N are one-potential

densities that share the stated properties for u , then

$$(1.21) \quad \prod_{j=1}^N \left(\frac{u_j(\bullet) + u_j(-\bullet)}{2} \right) \in L_{loc}^1(\mathbf{R}^d) \quad \Rightarrow \quad \prod_{j=1}^N \left(\frac{u_j(\bullet) + u_j(-\bullet)}{2} \right) \in L^1(\mathbf{R}^d).$$

If this is so, then the results of this paper imply Conjecture (1.20).

2. The stationary additive Lévy random field

Consider a classical Lévy process $X := \{X(t)\}_{t \geq 0}$ on \mathbf{R}^d with characteristic exponent Ψ . Let us introduce an independent copy X' of X , and extend the definition of X to a process indexed by \mathbf{R} as follows:

$$(2.1) \quad \tilde{X}(t) := \begin{cases} X(t) & \text{if } t \geq 0, \\ -X'(-t) & \text{if } t < 0. \end{cases}$$

This is the two-sided Lévy process with exponent Ψ in the sense that $\tilde{X} := \{\tilde{X}(t)\}_{t \in \mathbf{R}}$ has stationary and independent increments. Moreover, $\{\tilde{X}(t+s) - \tilde{X}(s)\}_{t \geq 0}$ is a copy of X for all $s \in \mathbf{R}$.

We also define $\tilde{\mathfrak{X}} := \{\tilde{\mathfrak{X}}(\mathbf{t})\}_{\mathbf{t} \in \mathbf{R}^N}$ as the corresponding N -parameter process, indexed by all of \mathbf{R}^N , whose values are in \mathbf{R}^d and are defined as

$$(2.2) \quad \tilde{\mathfrak{X}}(\mathbf{t}) := \tilde{X}_1(t_1) + \cdots + \tilde{X}_N(t_N) \quad \text{for all } \mathbf{t} := (t_1, \dots, t_N) \in \mathbf{R}^N.$$

We are assuming, of course, that $\tilde{X}_1, \dots, \tilde{X}_N$ are independent two-sided extensions of the processes X_1, \dots, X_N , respectively.

We intend to prove the following two-sided version of Theorem 1.1.

Theorem 2.1. *For all Borel sets $F \subseteq \mathbf{R}^d$,*

$$(2.3) \quad \mathbb{E} \left[\lambda_d \left(\tilde{\mathfrak{X}}(\mathbf{R}^N) \oplus F \right) \right] > 0 \quad \text{if and only if} \quad \text{cap}_{\Psi}(F) > 0.$$

This implies Theorem 1.1 effortlessly. Indeed, we know already from Remark 1.2 of Khoshnevisan, Xiao, and Zhong (2003) that

$$(2.4) \quad \text{cap}_{\Psi}(F) > 0 \quad \Rightarrow \quad \mathbb{E} \left[\lambda_d \left(\mathfrak{X}(\mathbf{R}_+^N) \oplus F \right) \right] > 0.$$

Thus, we seek only to derive the converse implication. But that follows from Theorem 2.1, because $\mathfrak{X}(\mathbf{R}_+^N) \subseteq \tilde{\mathfrak{X}}(\mathbf{R}^N)$.

Henceforth, we assume that the underlying probability space is the collection of all paths $\omega : \mathbf{R}^N \rightarrow \mathbf{R}^d$ that have the form $\omega(\mathbf{t}) = \sum_{j=1}^N \omega_j(t_j)$ for all $\mathbf{t} \in \mathbf{R}^N$, where each ω_j maps \mathbf{R}^N to \mathbf{R}^d such that $\omega_j(\mathbf{0}) = 0$; and $\omega_j \in D_{\mathbf{R}^d}(\mathbf{R})$, the Skorohod space of cadlag functions from \mathbf{R} —not $[0, \infty)$ —to \mathbf{R}^d .

We can then assume that the stationary additive Lévy fields, described earlier in this section, are in *canonical form*. That is, $\tilde{\mathfrak{X}}(\mathbf{t})(\omega) := \omega(\mathbf{t})$ for all $\mathbf{t} \in \mathbf{R}^N$ and $\omega \in \Omega$. Because we are interested only in distributional results, this is a harmless assumption.

Define P_x to be the law of $x + \tilde{\mathfrak{X}}$, and E_x the expectation operation with respect to P_x , for every $x \in \mathbf{R}^d$. Thus, we identify P with P_0 , and E with E_0 .

We are interested primarily in the σ -finite measure

$$(2.5) \quad P_{\lambda_d} := \int_{\mathbf{R}^d} P_x dx,$$

and the corresponding expectation operator E_{λ_d} , defined by

$$(2.6) \quad E_{\lambda_d} f := \int_{\Omega} f(\omega) P_{\lambda_d}(d\omega) \quad \text{for all } f \in L^1(P_{\lambda_d}).$$

If $A \ominus B := \{a - b : a \in A, b \in B\}$, then by the Fubini-Tonelli theorem,

$$(2.7) \quad \begin{aligned} E \left[\lambda_d \left(\tilde{\mathfrak{X}}(\mathbf{R}^N) \ominus F \right) \right] &= E \left[\int_{\mathbf{R}^d} \mathbf{1}_{\tilde{\mathfrak{X}}(\mathbf{R}^N) \ominus F}(x) dx \right] \\ &= \int_{\mathbf{R}^d} P_{-x} \left\{ \tilde{\mathfrak{X}}(\mathbf{t}) \in F \text{ for some } \mathbf{t} \in \mathbf{R}^N \right\} dx \\ &= P_{\lambda_d} \left\{ \tilde{\mathfrak{X}}(\mathbf{R}^N) \cap F \neq \emptyset \right\}. \end{aligned}$$

Thus, Theorem 2.1 is a potential-theoretic characterization of all polar sets for $\tilde{\mathfrak{X}}$ under the σ -finite measure P_{λ_d} . With this viewpoint in mind, we proceed to introduce some of the fundamental objects that are related to the process $\tilde{\mathfrak{X}}$.

Define, for all $\mathbf{t} \in \mathbf{R}^N$, the linear operator $P_{\mathbf{t}}$ as follows:

$$(2.8) \quad (P_{\mathbf{t}}f)(x) := E_x \left[f \left(\tilde{\mathfrak{X}}(\mathbf{t}) \right) \right] \quad \text{for all } x \in \mathbf{R}^d.$$

This is well-defined, for example, if $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$ is Borel-measurable, or if $f : \mathbf{R}^d \rightarrow \mathbf{R}$ is Borel-measurable and $P_{\mathbf{t}}(|f|)$ is finite at x . Also define the linear operator R by

$$(2.9) \quad (Rf)(x) := \frac{1}{2^N} \int_{\mathbf{R}^N} (P_{\mathbf{t}}f)(x) e^{-[\mathbf{t}]} d\mathbf{t} \quad \text{for all } x \in \mathbf{R}^d,$$

where

$$(2.10) \quad [\mathbf{t}] := |t_1| + \cdots + |t_N|$$

denotes the ℓ^1 -norm of $\mathbf{t} \in \mathbf{R}^N$. [We will use this notation throughout.] The ℓ^2 -norm of $\mathbf{t} \in \mathbf{R}^N$ will be denoted by $\|\mathbf{t}\|$.

Again, $(Rf)(x)$ is well defined if $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$ is Borel-measurable, or if $f : \mathbf{R}^d \rightarrow \mathbf{R}$ is Borel-measurable and $R(|f|)$ is finite at x .

Our next result is a basic regularity lemma for R . It should be recognized as a multiparameter version of a very well-known property of Markovian semigroups and their resolvents.

Lemma 2.2. *Each P_t and R are contractions on $L^p(\mathbf{R}^d)$, as long as $1 \leq p \leq \infty$.*

Proof. Choose and fix $j \in \{1, \dots, N\}$ and $t \in \mathbf{R}$, and define $\mu_{j,t}$ to be the distribution of the random variable $-\tilde{X}_j(t)$. If $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$ is Borel-measurable, then

$$(2.11) \quad P_t f = f * \mu_{1,t_1} * \dots * \mu_{N,t_N},$$

where $*$ denotes convolution. This implies readily that P_t is a contraction on $L^p(\mathbf{R}^d)$ for all $p \in [1, \infty]$, and hence,

$$(2.12) \quad \|Rf\|_{L^p(\mathbf{R}^d)} \leq \frac{1}{2^N} \int_{\mathbf{R}^N} \|P_t f\|_{L^p(\mathbf{R}^d)} e^{-[t]} dt.$$

Since P_t is a contraction on $L^p(\mathbf{R}^d)$, the preceding is bounded above by $\|f\|_{L^p(\mathbf{R}^d)}$. \square

Henceforth, let “ $\widehat{}$ ” denote the [Schwartz] Fourier transform on any and every Euclidean space \mathbf{R}^k . Our Fourier transform is normalized such that

$$(2.13) \quad \hat{f}(\xi) := \int_{\mathbf{R}^k} e^{i\xi \cdot x} f(x) dx,$$

for all $f \in L^1(\mathbf{R}^k)$, say. We then have the following.

Lemma 2.3. *If $f \in L^1(\mathbf{R}^d)$, then $\hat{R}f = K_\Psi$. That is, $(\widehat{Rf})(\xi) = K_\Psi(\xi) \hat{f}(\xi)$ for all $\xi \in \mathbf{R}^d$.*

Proof. Recall $\mu_{j,t}$ from the proof of Lemma 2.2. Its Fourier transform is given by

$$(2.14) \quad \widehat{\mu_{j,t}}(\xi) = \exp\{-|t|\Psi_j(-\text{sgn}(t)\xi)\} \quad \text{for all } \xi \in \mathbf{R}^d.$$

Equations (2.11) and (2.14), and the Plancherel theorem together imply that

$$(2.15) \quad (\widehat{P_t f})(\xi) = \hat{f}(\xi) \exp\left(-\sum_{j=1}^N |t_j| \Psi_j(-\text{sgn}(t_j)\xi)\right).$$

Consequently,

$$(2.16) \quad \begin{aligned} (\widehat{Rf})(\xi) &= \frac{1}{2^N} \hat{f}(\xi) \int_{\mathbf{R}^N} \exp\left(-\sum_{j=1}^N |t_j| [1 + \Psi_j(-\text{sgn}(t_j)\xi)]\right) dt \\ &= \frac{1}{2^N} \hat{f}(\xi) \prod_{j=1}^N \int_0^\infty (e^{-t[1+\Psi_j(\xi)]} + e^{-t[1+\Psi_j(-\xi)]}) dt. \end{aligned}$$

A direct computation reveals that $\widehat{Rf} = K_\Psi \hat{f}$, as asserted. \square

The following is a functional-analytic consequence.

Corollary 2.4. *The operator R maps $L^2(\mathbf{R}^d)$ into $L^2(\mathbf{R}^d)$, and is self-adjoint.*

Proof. By Lemma 2.3, if $f \in L^1(\mathbf{R}^d) \cap L^2(\mathbf{R}^d)$ and $g \in L^2(\mathbf{R}^d)$, then

$$(2.17) \quad \int_{\mathbf{R}^d} (Rf)(x)g(x) dx = \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} K_{\Psi}(\xi) \hat{f}(\xi) \overline{\hat{g}(\xi)} d\xi.$$

Thanks to Lemma 2.2, the preceding holds for all $f \in L^2(\mathbf{R}^d)$. Duality then implies that $R : L^2(\mathbf{R}^d) \rightarrow L^2(\mathbf{R}^d)$. Moreover, since K_{Ψ} is real, R is self-adjoint. \square

The following lemma shows that for every $\mathbf{t} \in \mathbf{R}^N$, the distribution of $\tilde{\mathfrak{X}}(\mathbf{t})$ under P_{λ} is λ_d . This is the reason why we call $\tilde{\mathfrak{X}}$ a *stationary* additive Lévy process.

Lemma 2.5. *If $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$ is Borel-measurable, then*

$$(2.18) \quad \mathbb{E}_{\lambda_d} \left[f \left(\tilde{\mathfrak{X}}(\mathbf{t}) \right) \right] = \int_{\mathbf{R}^d} (P_{\mathbf{t}}f)(x) dx = \int_{\mathbf{R}^d} f(y) dy \quad \text{for all } \mathbf{t} \in \mathbf{R}^N.$$

Proof. We apply the Fubini-Tonelli theorem to find that

$$(2.19) \quad \mathbb{E}_{\lambda_d} \left[f \left(\tilde{\mathfrak{X}}(\mathbf{t}) \right) \right] = \mathbb{E} \int_{\mathbf{R}^d} f \left(x + \tilde{\mathfrak{X}}(\mathbf{t}) \right) dx.$$

A change of variables $[y := x + \tilde{\mathfrak{X}}(\mathbf{t})]$ proves that the preceding expression is equal to the integral of f . This implies half of the lemma. Another application of the Fubini-Tonelli theorem implies the remaining half as well. \square

Let us choose and fix $\pi \subseteq \{1, \dots, N\}$ and identify π with the partial order \prec_{π} , on \mathbf{R}^N , which is defined as follows: For all $\mathbf{s}, \mathbf{t} \in \mathbf{R}^N$,

$$(2.20) \quad \mathbf{s} \prec_{\pi} \mathbf{t} \quad \text{iff} \quad \begin{cases} s_i \leq t_i & \text{for all } i \in \pi, \text{ and} \\ s_i > t_i & \text{for all } i \notin \pi. \end{cases}$$

The collection of all $\pi \subseteq \{1, \dots, N\}$ forms a collective total order on \mathbf{R}^N in the sense that

$$(2.21) \quad \text{for all } \mathbf{s}, \mathbf{t} \in \mathbf{R}^N \text{ there exists } \pi \subseteq \{1, \dots, N\} \text{ such that } \mathbf{s} \prec_{\pi} \mathbf{t}.$$

For all $\pi \subseteq \{1, \dots, N\}$, we define the π -*history* of the random field $\tilde{\mathfrak{X}}$ as the collection

$$(2.22) \quad \mathcal{H}_{\pi}(\mathbf{t}) := \sigma \left(\left\{ \tilde{\mathfrak{X}}(\mathbf{s}) \right\}_{\mathbf{s} \prec_{\pi} \mathbf{t}} \right) \quad \text{for all } \mathbf{t} \in \mathbf{R}^N,$$

where $\sigma(\dots)$ denotes the σ -algebra generated by whatever is in the parentheses. Without loss of generality, we assume that each $\mathcal{H}_{\pi}(\mathbf{t})$ is complete with respect to P_x for all $x \in \mathbf{R}^d$; else, we replace it with the said completion. Also, we assume without loss of generality that $\mathbf{t} \mapsto \mathcal{H}_{\pi}(\mathbf{t})$ is π -right-continuous. More precisely, we assume that

$$(2.23) \quad \mathcal{H}_{\pi}(\mathbf{t}) = \bigcap_{\mathbf{s} \in \mathbf{R}^N : \mathbf{t} \prec_{\pi} \mathbf{s}} \mathcal{H}_{\pi}(\mathbf{s}) \quad \text{for all } \mathbf{t} \in \mathbf{R}^N \text{ and } \pi \subseteq \{1, \dots, N\}.$$

If not, then we replace the left-hand side by the right-hand side everywhere.

Proposition 2.6 (A Markov-random-field property). *Suppose $\pi \subseteq \{1, \dots, N\}$ and $\mathbf{s} \prec_\pi \mathbf{t}$, both in \mathbf{R}^N . Then, for all measurable functions $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$,*

$$(2.24) \quad \mathbb{E}_{\lambda_d} \left[f \left(\tilde{\mathbf{x}}(\mathbf{t}) \right) \mid \mathcal{H}_\pi(\mathbf{s}) \right] = (P_{\mathbf{t}-\mathbf{s}} f) \left(\tilde{\mathbf{x}}(\mathbf{s}) \right) \quad \mathbb{P}_{\lambda_d}\text{-a.s.}$$

Proof. Choose and fix Borel measurable functions $g, \phi_1, \dots, \phi_m : \mathbf{R}^d \rightarrow \mathbf{R}_+$, and “ N -parameter time points” $\mathbf{s}^1, \dots, \mathbf{s}^m \in \mathbf{R}^N$ such that

$$(2.25) \quad \mathbf{s}^j \prec_\pi \mathbf{s} \prec_\pi \mathbf{t} \quad \text{for all } j = 1, \dots, m.$$

According to the Fubini-Tonelli theorem,

$$(2.26) \quad \begin{aligned} & \mathbb{E}_{\lambda_d} \left[f \left(\tilde{\mathbf{x}}(\mathbf{t}) \right) g \left(\tilde{\mathbf{x}}(\mathbf{s}) \right) \prod_{j=1}^m \phi_j \left(\tilde{\mathbf{x}}(\mathbf{s}^j) \right) \right] \\ &= \int_{\mathbf{R}^d} \mathbb{E} \left[f \left(x + \tilde{\mathbf{x}}(\mathbf{t}) \right) g \left(x + \tilde{\mathbf{x}}(\mathbf{s}) \right) \prod_{j=1}^m \phi_j \left(x + \tilde{\mathbf{x}}(\mathbf{s}^j) \right) \right] dx \\ &= \int_{\mathbf{R}^d} \mathbb{E} \left[f(A + y) \prod_{j=1}^m \phi_j(A_j + y) \right] g(y) dy, \end{aligned}$$

where $A := \tilde{\mathbf{x}}(\mathbf{t}) - \tilde{\mathbf{x}}(\mathbf{s})$ and $A_j := \tilde{\mathbf{x}}(\mathbf{s}^j) - \tilde{\mathbf{x}}(\mathbf{s})$ for all $j = 1, \dots, m$.

The independent-increments property of each of the Lévy processes \tilde{X}_j implies that A is independent of $\{A_j\}_{j=1}^m$. This and the stationary-increments property of $\tilde{X}_1, \dots, \tilde{X}_N$ together imply that

$$(2.27) \quad \begin{aligned} & \mathbb{E}_{\lambda_d} \left[f \left(\tilde{\mathbf{x}}(\mathbf{t}) \right) g \left(\tilde{\mathbf{x}}(\mathbf{s}) \right) \prod_{j=1}^m \phi_j \left(\tilde{\mathbf{x}}(\mathbf{s}^j) \right) \right] \\ &= \int_{\mathbf{R}^d} \mathbb{E} [f(A + y)] \mathbb{E} \left[\prod_{j=1}^m \phi_j(A_j + y) \right] g(y) dy. \end{aligned}$$

After a change of variables and an appeal to the stationary-independent property of the increments of X_1, \dots, X_N and X'_1, \dots, X'_N , we arrive at the following:

$$\begin{aligned}
 (2.28) \quad & \mathbb{E}_{\lambda_d} \left[f(\tilde{\mathbf{x}}(\mathbf{t})) g(\tilde{\mathbf{x}}(\mathbf{s})) \prod_{j=1}^m \phi_j(\tilde{\mathbf{x}}(\mathbf{s}^j)) \right] \\
 &= \int_{\mathbf{R}^d} \mathbb{E} \left[f(\tilde{\mathbf{x}}(\mathbf{t} - \mathbf{s}) + y) \right] \mathbb{E}_y \left[\prod_{j=1}^m \phi_j(A_j) \right] g(y) dy \\
 &= \int_{\mathbf{R}^d} (P_{\mathbf{t}-\mathbf{s}}f)(y) \mathbb{E}_y \left[\prod_{j=1}^m \phi_j(A_j) \right] g(y) dy.
 \end{aligned}$$

By the monotone class theorem, for all nonnegative $\mathcal{H}_\pi(\mathbf{s})$ -measurable random variables Y ,

$$(2.29) \quad \mathbb{E}_{\lambda_d} \left[f(\tilde{\mathbf{x}}(\mathbf{t})) g(\tilde{\mathbf{x}}(\mathbf{s})) Y \right] = \int_{\mathbf{R}^d} (P_{\mathbf{t}-\mathbf{s}}f)(y) \psi(y) g(y) dy,$$

where $\psi : \mathbf{R}^d \rightarrow \mathbf{R}_+$ is a measurable function. This proves the proposition. \square

Lemma 2.7. *If $f, g \in L^2(\mathbf{R}^d)$ and $\mathbf{t}, \mathbf{s} \in \mathbf{R}^N$, then*

$$(2.30) \quad \mathbb{E}_{\lambda_d} \left[f(\tilde{\mathbf{x}}(\mathbf{s})) g(\tilde{\mathbf{x}}(\mathbf{t})) \right] = \int_{\mathbf{R}^d} f(y) (P_{\mathbf{t}-\mathbf{s}}g)(y) dy.$$

Proof. We may consider, without loss of generality, measurable and nonnegative functions $f, g \in L^2(\mathbf{R}^d)$. Let π denote the collection of all $i \in \{1, \dots, N\}$ such that $s_i \leq t_i$. Then $\mathbf{s} \prec_\pi \mathbf{t}$, and Proposition 2.6 implies that \mathbb{P}_{λ_d} -a.s.,

$$(2.31) \quad \mathbb{E}_{\lambda_d} \left[f(\tilde{\mathbf{x}}(\mathbf{t})) \mid \mathcal{H}_\pi(\mathbf{s}) \right] = (P_{\mathbf{t}-\mathbf{s}}f)(\tilde{\mathbf{x}}(\mathbf{s})).$$

This and Lemma 2.5 together conclude the proof. \square

3. The sojourn operator

Recall (2.10), and consider the “sojourn operator,”

$$(3.1) \quad Sf := \frac{1}{2^N} \int_{\mathbf{R}^N} f(\tilde{\mathbf{x}}(\mathbf{t})) e^{-[\mathbf{t}]} d\mathbf{t}.$$

Our first lemma records the fact that S maps density functions to mean-one random variables $[\mathbb{P}_{\lambda_d}]$.

Lemma 3.1. *If f is a probability density function on \mathbf{R}^d , then $\mathbb{E}_{\lambda_d}[Sf] = 1$.*

This follows readily from Lemma 2.5. Our next result shows that, under a mild condition on Ψ_j , S embeds functions in $L^2(\mathbf{R}^d)$, quasi-isometrically, into the subcollection of all functions in $L^2(\mathbf{R}^d)$ that have finite energy. Namely,

Proposition 3.2. *If $f \in L^2(\mathbf{R}^d)$, then*

$$(3.2) \quad \|Sf\|_{L^2(\mathbf{P}_{\lambda_d})} \leq \sqrt{I_{\Psi}(f)}.$$

Suppose, in addition, that there exists a constant $c \in (0, \sqrt{2})$, such that the following sector condition holds for all $j = 1, \dots, N$:

$$(3.3) \quad |\operatorname{Im} \Psi_j(\xi)| \leq c(1 + \operatorname{Re} \Psi_j(\xi)) \quad \text{for all } \xi \in \mathbf{R}^d.$$

Then, there exists a constant $A \in (0, 1)$ such that

$$(3.4) \quad A \sqrt{I_{\Psi}(f)} \leq \|Sf\|_{L^2(\mathbf{P}_{\lambda_d})} \leq \sqrt{I_{\Psi}(f)}.$$

Remark 3.3 (Generalized Sobolev spaces). When $N = 1$ and the Lévy process in question is symmetric, the following problem arises in the theory of Dirichlet forms: For what f in the class $\mathcal{D}(\mathbf{R}^d)$, of Schwartz distributions on \mathbf{R}^d , can we define Sf as an element of $L^2(\mathbf{P}_{\lambda_d})$ (say)? This problem continues to make sense in the more general context of additive Lévy processes. And the answer is given by (3.4) in Proposition 3.2 as follows: Assume that the *sector condition* (3.3) holds for all $j = 1, \dots, N$ and some $c \in (0, \sqrt{2})$. Let $\mathcal{S}_{\Psi}(\mathbf{R}^d)$ denote the completion of the collection of all members of $L^2(\mathbf{R}^d)$ that have finite energy I_{Ψ} , where the completion is made in the norm $\|f\|_{\Psi} := I_{\Psi}^{1/2}(f) + \|f\|_{L^2(\mathbf{R}^d)}$. Then, there exists an a.s.-unique maximal extension \bar{S} of S such that $\bar{S} : \mathcal{S}_{\Psi}(\mathbf{R}^d) \rightarrow \bar{S}(\mathcal{S}_{\Psi}(\mathbf{R}^d))$ is a quasi-isometry. The space $\mathcal{S}_{\Psi}(\mathbf{R}^d)$ generalizes further some of the ψ -Bessel potential spaces of Farkas, Jacob, and Schilling (2001) and Farkas and Leopold (2006); see also Jacob and Schilling (2005), Masja and Nagel (1978), and Slobodeckii (1958). \square

The proof requires a technical lemma, which we develop first.

Lemma 3.4. *For all $z \in \mathbf{C}$ define*

$$(3.5) \quad \Lambda(z) := \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-|t|-|s|-|t-s|\sigma(z;t-s)} dt ds,$$

where $\sigma(z; r) := z$ if $r \geq 0$ and $\sigma(z; r) := \bar{z}$ otherwise. Then for all $z \in \mathbf{C}$ with $\operatorname{Re} z \geq 0$,

$$(3.6) \quad \Lambda(z) \leq 4 \operatorname{Re} \left(\frac{1}{1+z} \right).$$

If, in addition, $|\operatorname{Im} z| \leq c(1 + \operatorname{Re} z)$ for some $c \in (0, \sqrt{2})$, then

$$(3.7) \quad \Lambda(z) \geq 2(2 - c^2) \operatorname{Re} \left(\frac{1}{1+z} \right).$$

Proof. The double integral is computed by dividing the region of integration into four natural parts: (i) $s, t \geq 0$; (ii) $s, t \leq 0$; (iii) $t \geq 0 \geq s$; and (iv) $s \geq 0 \geq t$. Direct computation

reveals that for all $z \in \mathbf{C}$ with $\operatorname{Re} z \geq 0$

$$(3.8) \quad \int_0^\infty \int_0^\infty e^{-|t|-|s|-|t-s|\sigma(z;t-s)} dt ds + \int_{-\infty}^0 \int_{-\infty}^0 e^{-|t|-|s|-|t-s|\sigma(z;t-s)} dt ds = 2 \operatorname{Re} \left(\frac{1}{1+z} \right).$$

Similarly, one can compute

$$(3.9) \quad \int_0^\infty \int_{-\infty}^0 e^{-|t|-|s|-|t-s|\sigma(z;t-s)} dt ds + \int_{-\infty}^0 \int_0^\infty e^{-|t|-|s|-|t-s|\sigma(z;t-s)} dt ds = \frac{1}{(1+z)^2} + \frac{1}{(1+\bar{z})^2}.$$

Consequently,

$$(3.10) \quad \Lambda(z) = 2 \operatorname{Re} \left(\frac{1}{1+z} \right) + \frac{2((1+\operatorname{Re} z)^2 - (\operatorname{Im} z)^2)}{|1+z|^4},$$

for all $z \in \mathbf{C}$ with $\operatorname{Re} z \geq 0$. It follows that

$$(3.11) \quad \Lambda(z) \leq 2 \operatorname{Re} \left(\frac{1}{1+z} \right) \left[1 + \operatorname{Re} \left(\frac{1}{1+z} \right) \right]$$

for all $z \in \mathbf{C}$ with $\operatorname{Re} z \geq 0$. Whenever $\operatorname{Re} z \geq 0$, we have $0 \leq \operatorname{Re}(1+z)^{-1} \leq 1$, and hence (3.6) follows from (3.11). On the other hand, if $|\operatorname{Im} z| \leq c(1+\operatorname{Re} z)$, then (3.10) yields

$$(3.12) \quad \Lambda(z) \geq 2 \operatorname{Re} \left(\frac{1}{1+z} \right) + 2(1-c^2) \left[\operatorname{Re} \left(\frac{1}{1+z} \right) \right]^2,$$

from which the result follows readily, because $0 \leq \operatorname{Re}[(1+z)^{-1}] \leq 1$ when $\operatorname{Re} z \geq 0$. \square

Proof of Proposition 3.2. We apply Lemma 2.7 to deduce that

$$(3.13) \quad \mathbb{E}_{\lambda_d}(|Sf|^2) = \frac{1}{4^N} \int_{\mathbf{R}^d} \iint_{\mathbf{R}^N \times \mathbf{R}^N} e^{-[t]-[s]} f(y) (P_{t-s}f)(y) dt ds dy.$$

In accord with (2.15) and Parseval's identity, for all $\mathbf{u} \in \mathbf{R}^N$,

$$(3.14) \quad \begin{aligned} \int_{\mathbf{R}^d} f(y) (P_{\mathbf{u}}f)(y) dy &= \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \overline{\hat{f}(\xi)} \widehat{(P_{\mathbf{u}}f)}(\xi) d\xi \\ &= \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} |\hat{f}(\xi)|^2 \exp \left(- \sum_{j=1}^N |u_j| \Psi_j(-\operatorname{sgn}(u_j)\xi) \right) d\xi. \end{aligned}$$

This and the Fubini-Tonelli theorem together reveal that

$$(3.15) \quad \mathbb{E}_{\lambda_d}(|Sf|^2) = \frac{1}{4^N (2\pi)^d} \int_{\mathbf{R}^d} |\hat{f}(\xi)|^2 \prod_{j=1}^N \Lambda(\Psi_j(\xi)) d\xi.$$

Since $\operatorname{Re} \Psi_j(\xi) \geq 0$ for all $\xi \in \mathbf{R}^d$ and $j = 1, \dots, N$, we apply Lemma 3.4 to this formula, and conclude the proof of the proposition. \square

4. Proof of Theorem 2.1

Thanks to the definition of $\operatorname{cap}_{\Psi}$, and to the countable additivity of P , it suffices to consider only the case that

$$(4.1) \quad F \text{ is a compact set.}$$

This condition is tacitly assumed throughout this section. We note, in particular, that $\mathcal{P}_c(F)$ denotes merely the collection of all Borel probability measures that are supported on F .

Proposition 5.7 of Khoshnevisan, Xiao, and Zhong (2003) proves that for every compact set $F \subseteq \mathbf{R}^d$,

$$(4.2) \quad \operatorname{cap}_{\Psi}(F) > 0 \implies \mathbb{E} [\lambda_d (\mathfrak{X}([0, r]^N) \oplus F)] > 0$$

for all $r > 0$. It is clear the latter implies that $\mathbb{E} [\lambda_d (\tilde{\mathfrak{X}}(\mathbf{R}^N) \oplus F)] > 0$, and one obtains half of the theorem.

Since $\operatorname{cap}_{\Psi}(-F) = \operatorname{cap}_{\Psi}(F)$, we may and will replace F by $-F$ throughout. In light of (2.7), we then assume that

$$(4.3) \quad P_{\lambda_d} \left\{ \tilde{\mathfrak{X}}(\mathbf{R}^N) \cap F \neq \emptyset \right\} > 0,$$

and seek to deduce the existence of a probability measure μ on F such that $I_{\Psi}(\mu) < \infty$.

We will prove a little more. Namely, that for all $k > 0$ there exists a constant $A = A(k, N) \in (0, \infty)$ such that

$$(4.4) \quad P_{\lambda_d} \left\{ \tilde{\mathfrak{X}}([-k, k]^N) \cap F \neq \emptyset \right\} \leq A \operatorname{cap}_{\Psi}(F).$$

In fact, we will prove that for all sufficiently large $k > 0$,

$$(4.5) \quad P_{\lambda_d} \left\{ \tilde{\mathfrak{X}}([-k, k]^N) \cap F \neq \emptyset \right\} \leq e^{2Nk} 4^N \operatorname{cap}_{\Psi}(F).$$

This would conclude our proof of the second half of the theorem.

It is a standard fact that there exists a probability density function ϕ_1 in $C^\infty(\mathbf{R}^d)$ with the following properties:

- P1.** $\phi_1(x) = 0$ if $\|x\| > 1$;
- P2.** $\phi_1(x) = \phi_1(-x)$ for all $x \in \mathbf{R}^d$;
- P3.** $\hat{\phi}_1(x) \geq 0$ for all $x \in \mathbf{R}^d$.

This can be obtained, for example, readily from Plancherel's [duality] theorem of Fourier analysis.

We recall also the following standard fact: $\phi_1 \in L^1(\mathbf{R}^d)$ and **P3** together imply that $\hat{\phi}_1 \in L^1(\mathbf{R}^d)$ (Hawkes, 1984, Lemma 1).

Now we define an approximation to the identity $\{\phi_\epsilon\}_{\epsilon>0}$ by setting

$$(4.6) \quad \phi_\epsilon(x) := \frac{1}{\epsilon^d} \phi_1\left(\frac{x}{\epsilon}\right) \quad \text{for all } x \in \mathbf{R}^d \text{ and } \epsilon > 0.$$

It follows readily from this that for every $\mu \in \mathcal{P}_c(F)$:

- (i) $\mu * \phi_\epsilon$ is a uniformly continuous probability density;
- (ii) $\mu * \phi_\epsilon$ is supported on the closed ϵ -enlargement of F , which we denote by F_ϵ , for all $\epsilon > 0$;
- (iii) ϕ_ϵ is symmetric; and
- (iv) $\lim_{\epsilon \rightarrow 0^+} \hat{\phi}_\epsilon(\xi) = 1$ for all $\xi \in \mathbf{R}^d$.

As was done in Khoshnevisan, Xiao, and Zhong (2003), we can find a random variable \mathbf{T} with values in $\mathbf{R}^N \cup \{\infty\}$ such that:

- (1) $\{\mathbf{T} = \infty\}$ is equal to the event that $\tilde{\mathbf{X}}(t) \notin F$ for all $t \in \mathbf{R}^N$;
- (2) $\tilde{\mathbf{X}}(\mathbf{T}) \in F$ on $\{\mathbf{T} \neq \infty\}$.

This can be accomplished pathwise. Consequently, (4.3) is equivalent to the condition that

$$(4.7) \quad \mathbb{P}_{\lambda_d} \{\mathbf{T} \neq \infty\} > 0.$$

Define for all Borel sets $A \subseteq \mathbf{R}^d$ and all integers $k \geq 1$,

$$(4.8) \quad \mu_k(A) := \mathbb{P}_{\lambda_d} \left(\tilde{\mathbf{X}}(\mathbf{T}) \in A \mid \mathbf{T} \in [-k, k]^N \right).$$

We claim that μ_k is a probability measure on F for all k sufficiently large. In order to prove this claim we choose and fix $l > 0$, and consider

$$(4.9) \quad \mu_{k,l}(A) := \frac{\mathbb{P}_{\lambda_d} \left\{ \tilde{\mathbf{X}}(\mathbf{T}) \in A, \mathbf{T} \in [-k, k]^N, \left| \tilde{\mathbf{X}}(\mathbf{0}) \right| \leq l \right\}}{\mathbb{P}_{\lambda_d} \left\{ \mathbf{T} \in [-k, k]^N, \left| \tilde{\mathbf{X}}(\mathbf{0}) \right| \leq l \right\}}.$$

Because the \mathbb{P}_{λ_d} -distribution of $\tilde{\mathbf{X}}(\mathbf{0})$ is λ_d , $\mu_{k,l}$ is a probability measure on F for all k and l sufficiently large; see (4.7). And $\mu_{k,l}(A)$ converges to $\mu_k(A)$ for all Borel sets A as $l \uparrow \infty$ [monotone convergence theorem]. This proves the assertion that $\mu_k \in \mathcal{P}_c(F)$ for all k large. Choose and fix such a large integer k .

Now we define $f_\epsilon := \mu_k * \phi_\epsilon$, and observe that according to Proposition 2.6, for all non-random times $\tau \in \mathbf{R}^N$,

$$\begin{aligned}
(4.10) \quad & \sum_{\pi \subseteq \{1, \dots, N\}} \mathbb{E}_{\lambda_d} \left[\int_{t \succ_\pi \tau} f_\epsilon(\tilde{\mathbf{x}}(t)) e^{-[t]} dt \mid \mathcal{H}_\pi(\tau) \right] \\
&= \sum_{\pi \subseteq \{1, \dots, N\}} \int_{t \succ_\pi \tau} (P_{t-\tau} f_\epsilon)(\tilde{\mathbf{x}}(\tau)) e^{-[t]} dt \\
&\geq e^{-[\tau]} \sum_{\pi \subseteq \{1, \dots, N\}} \int_{s \succ_\pi \mathbf{0}} (P_s f_\epsilon)(\tilde{\mathbf{x}}(\tau)) e^{-[s]} ds.
\end{aligned}$$

We can apply (2.21) to deduce then that for all nonrandom times $\tau \in \mathbf{R}^N$,

$$\begin{aligned}
(4.11) \quad & \mathbb{E}_{\lambda_d} [Sf_\epsilon \mid \mathcal{H}_\pi(\tau)] \geq \frac{e^{-[\tau]}}{2^N} \int_{\mathbf{R}^N} (P_s f_\epsilon)(\tilde{\mathbf{x}}(\tau)) e^{-[s]} ds \\
&= e^{-[\tau]} (Rf_\epsilon)(\tilde{\mathbf{x}}(\tau)).
\end{aligned}$$

Because f_ϵ is continuous and compactly supported, one can verify from (2.9) that Rf_ϵ is continuous [this can also be shown by Lemma 2.3, the fact that $\hat{f}_\epsilon \in L^1(\mathbf{R}^d)$ and the Fourier inversion formula]. Also, the $L^2(\mathbf{P}_{\lambda_d})$ -norm of the left-most term in (4.11) is bounded above by $\|Sf_\epsilon\|_{L^2(\mathbf{P}_{\lambda_d})}$, and this is at most $\sqrt{I_\Psi(f_\epsilon)}$, in turn; see Proposition 3.2. It is easy to see that

$$\begin{aligned}
(4.12) \quad & I_\Psi(f_\epsilon) = I_\Psi(\mu_k * \phi_\epsilon) \\
&\leq \frac{1}{(2\pi)^{2d}} \int_{\mathbf{R}^d} |\hat{\phi}_\epsilon(\xi)|^2 d\xi \\
&< \infty.
\end{aligned}$$

Since $|\hat{\mu}_k(\xi)|^2 K_\Psi(\xi) \leq 1$, we conclude that $\|Sf_\epsilon\|_{L^2(\mathbf{P}_{\lambda_d})}$ is finite for all $\epsilon > 0$.

A simple adaptation of the proof of Proposition 2.6 proves that \mathcal{H}_π is a *commuting filtration* for all $\pi \subseteq \{1, \dots, N\}$. By this we mean that for all $t \in \mathbf{R}$,

$$(4.13) \quad \mathcal{H}_\pi^1(t_1), \dots, \mathcal{H}_\pi^N(t_N) \text{ are conditionally independent } [\mathbf{P}_{\lambda_d}], \text{ given } \mathcal{H}_\pi(t),$$

where $\mathcal{H}_\pi^j(t_j)$ is defined as the following σ -algebra:

$$(4.14) \quad \mathcal{H}_\pi^j(t_j) := \begin{cases} \sigma(\tilde{X}_j(s); -\infty < s \leq t_j) & \text{if } j \in \pi, \\ \sigma(\tilde{X}_j(s); \infty > s \geq t_j) & \text{if } j \notin \pi. \end{cases}$$

Here, $\sigma(\dots)$ denotes the σ -algebra generated by the random variables in the parenthesis.

The stated commutation property readily implies that for all random variables $Y \in L^2(\mathbf{P}_{\lambda_d})$ and partial orders $\pi \subseteq \{1, \dots, N\}$:

- (1) $\boldsymbol{\tau} \mapsto \mathbb{E}_{\lambda_d}[Y \mid \mathcal{H}_\pi(\boldsymbol{\tau})]$ has a version that is cadlag in each of its N variables, uniformly in all other $N - 1$ variables; and
- (2) The second moment of $\sup_{\boldsymbol{\tau} \in \mathbf{R}^N} \mathbb{E}_{\lambda_d}[Y \mid \mathcal{H}_\pi(\boldsymbol{\tau})]$ is at most 4^N times the second moment of Y $[\mathbb{P}_{\lambda_d}]$.

See Khoshnevisan (2002, Theorem 2.3.2, p. 235) for the case where \mathbb{P}_{λ_d} is replaced by a probability measure. The details of the remaining changes are explained in a slightly different setting in Lemma 4.2 of Khoshnevisan, Xiao, and Zhong (2003). In summary, (4.11) holds for all $\boldsymbol{\tau} \in \mathbf{R}^N$, \mathbb{P}_{λ_d} -almost surely [note the order of the quantifiers]. It follows immediately from this that for all integers $k \geq 1$,

$$(4.15) \quad \sup_{\boldsymbol{\tau} \in \mathbf{R}^N} \mathbb{E}_{\lambda_d}[Sf_\epsilon \mid \mathcal{H}_\pi(\boldsymbol{\tau})] \geq e^{-Nk} (Rf_\epsilon) \left(\tilde{\mathbf{x}}(\mathbf{T}) \right) \cdot \mathbf{1}_{\{\mathbf{T} \in [-k, k]^N\}} \quad \mathbb{P}_{\lambda_d}\text{-a.s.}$$

According to Item (2) above, the second moment of the left-hand side is at most 4^N times the second moment of Sf_ϵ . As we noticed earlier, the latter is at most $I_\Psi(f_\epsilon)$. Therefore,

$$(4.16) \quad \begin{aligned} e^{2Nk} 4^N I_\Psi(f_\epsilon) &\geq \mathbb{E}_{\lambda_d} \left[\left| (Rf_\epsilon) \left(\tilde{\mathbf{x}}(\mathbf{T}) \right) \right|^2 ; \mathbf{T} \in [-k, k]^N \right] \\ &= \mathbb{E}_{\lambda_d} \left[\left| (Rf_\epsilon) \left(\tilde{\mathbf{x}}(\mathbf{T}) \right) \right|^2 \mid \mathbf{T} \in [-k, k]^N \right] \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \}. \end{aligned}$$

It follows from this and the Cauchy–Schwarz inequality that

$$(4.17) \quad e^{2Nk} 4^N I_\Psi(f_\epsilon) \geq \left| \mathbb{E}_{\lambda_d} \left[(Rf_\epsilon) \left(\tilde{\mathbf{x}}(\mathbf{T}) \right) \mid \mathbf{T} \in [-k, k]^N \right] \right|^2 \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \}.$$

Using the definition of μ_k , we can write the above as

$$(4.18) \quad \begin{aligned} e^{2Nk} 4^N I_\Psi(f_\epsilon) &\geq \left| \int_{\mathbf{R}^d} (Rf_\epsilon) d\mu_k \right|^2 \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \} \\ &= \frac{1}{(2\pi)^{2d}} \left| \int_{\mathbf{R}^d} \overline{\hat{\mu}_k(\xi)} \widehat{(Rf_\epsilon)}(\xi) d\xi \right|^2 \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \}, \end{aligned}$$

where the equality follows from the Parseval identity. According to Lemma 2.3, the Fourier transform of Rf_ϵ is K_Ψ times the Fourier transform of f_ϵ , and the latter is $\hat{\mu}_k \hat{\phi}_\epsilon$. Hence (4.18) implies that

$$(4.19) \quad e^{2Nk} 4^N I_\Psi(f_\epsilon) \geq \frac{1}{(2\pi)^{2d}} \left| \int_{\mathbf{R}^d} |\hat{\mu}_k(\xi)|^2 \hat{\phi}_\epsilon(\xi) K_\Psi(\xi) d\xi \right|^2 \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \}.$$

Now we apply Property **P3** to deduce that $\hat{\phi}_\epsilon \geq 0$. Because ϕ_ϵ is also a probability density, it follows that $\hat{\phi}_\epsilon \geq |\hat{\phi}_\epsilon|^2$. Consequently,

$$(4.20) \quad \begin{aligned} e^{2Nk} 4^N I_\Psi(f_\epsilon) &\geq \frac{1}{(2\pi)^{2d}} \left| \int_{\mathbf{R}^d} |\hat{f}_\epsilon(\xi)|^2 K_\Psi(\xi) d\xi \right|^2 \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \} \\ &= |I_\Psi(f_\epsilon)|^2 \cdot \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \}. \end{aligned}$$

We have seen in (4.12) that $I_\Psi(f_\epsilon)$ is finite for each $\epsilon > 0$. If it were zero for arbitrary small $\epsilon > 0$, then we apply Fatou's lemma to deduce that $I_\Psi(\mu_k) \leq \liminf_{\epsilon \rightarrow 0} I_\Psi(f_\epsilon) = 0$. This and the fact that $K_\Psi(\xi) > 0$ for all $\xi \in \mathbf{R}^d$ would imply $\hat{\mu}_k(\xi) \equiv 0$ for all $\xi \in \mathbf{R}^d$, which is a contradiction. Hence we can deduce that for all $\epsilon > 0$ small enough,

$$(4.21) \quad \frac{e^{2Nk} 4^N}{I_\Psi(f_\epsilon)} \geq \mathbb{P}_{\lambda_d} \{ \mathbf{T} \in [-k, k]^N \},$$

and this is positive for k large; see (4.7). The right-hand side of (4.21) is independent of $\epsilon > 0$. Hence, we can let $\epsilon \downarrow 0$ and appeal to Fatou's lemma to deduce that $I_\Psi(\mu_k) < \infty$. Thus, in any event, we have produced a probability measure μ_k on F whose energy $I_\Psi(\mu_k)$ is finite. This concludes the proof, and also implies (4.5), thanks to the defining properties of the function \mathbf{T} . \square

5. On kernels of positive type

In this section we study kernels of positive type; they are recalled next. Here and throughout, $\bar{\mathbf{R}}_+ := [0, \infty]$ is defined to be the usual one-point compactification of $\mathbf{R}_+ := [0, \infty)$, and is endowed with the corresponding Borel sigma-algebra.

Definition 5.1. A *kernel* [on \mathbf{R}^d] is a Borel measurable function $\kappa : \mathbf{R}^d \rightarrow \bar{\mathbf{R}}_+$ such that $\kappa \in L^1_{loc}(\mathbf{R}^d)$. If, in addition, $\hat{\kappa}(\xi) \geq 0$ for all $\xi \in \mathbf{R}^d$, then we say that κ is a kernel of *positive type*.

Clearly, every kernel κ can be redefined on a Lebesgue-null set so that the resulting modification $\tilde{\kappa}$ maps \mathbf{R}^d into \mathbf{R}_+ . However, we might lose some of the nice properties of κ by doing this. A notable property is that κ might be continuous; that is, $\kappa(x) \rightarrow \kappa(y)$ —in $\bar{\mathbf{R}}_+$ —as x converges to y in \mathbf{R}^d . In this case, $\tilde{\kappa}$ might not be continuous. From this perspective, it is sometimes advantageous to work with the $\bar{\mathbf{R}}_+$ -valued function κ . For examples, we have in mind *Riesz kernels*. They are defined as follows: Choose and fix some number $\alpha \in (0, d)$, and then let the Riesz kernel κ_α of index α be

$$(5.1) \quad \kappa_\alpha(x) := \begin{cases} \|x\|^{\alpha-d} & \text{if } x \neq 0, \\ \infty & \text{if } x = 0. \end{cases}$$

It is easy to check that κ_α is a continuous kernel for each $\alpha \in (0, d)$. In fact, every κ_α is a kernel of positive type, as can be seen via the following standard fact:

$$(5.2) \quad \hat{\kappa}_\alpha(\xi) = c_{d,\alpha} \kappa_{d-\alpha}(\xi) \quad \text{for all } \xi \in \mathbf{R}^d.$$

Here $c_{d,\alpha}$ is a universal constant that depends only on d and α ; see Kahane (1985, p. 134) or Mattila (1995, eq. (12.10), p. 161), for example.

It is true—but still harder to prove—that for all Borel probability measures μ on \mathbf{R}^d ,

$$(5.3) \quad \iint \frac{\mu(dx) \mu(dy)}{\|x - y\|^{d-\alpha}} = \frac{c_{d,\alpha}}{(2\pi)^d} \int_{\mathbf{R}^d} |\hat{\mu}(\xi)|^2 \frac{d\xi}{\|\xi\|^\alpha}.$$

See Mattila (1995, p. Lemma 12.12, p. 162).

The utility of (5.3) is in the fact that it shows that Riesz-type energies of the left-hand side are equal to Pólya–Szegő energies of the right-hand side. This is a probabilistically significant fact. For example, consider the case that $\alpha \in (0, 2]$. Then, $\mu \mapsto \iint \|x - y\|^{-d+\alpha} \mu(dx) \mu(dy)$ is the “energy functional” associated to continuous additive functionals of various stable processes of index α . At the same time, $\mu \mapsto (2\pi)^{-d} \int_{\mathbf{R}^d} |\hat{\mu}(\xi)|^2 \|\xi\|^{-\alpha} d\xi$ is a Fourier-analytic energy form of the type that appears more generally in the earlier parts of the present paper. Roughly speaking, $(2\pi)^{-d} \int_{\mathbf{R}^d} |\hat{\mu}(\xi)|^2 \|\xi\|^{-\alpha} d\xi \asymp I_\Psi(\mu)$, where $\Psi(\xi) = \|\xi\|^\alpha$ defines the Lévy exponent of an isotropic stable process of index α . [Analytically speaking, this is the Sobolev norm of μ that corresponds to the fractional Laplacian operator $-(-\Delta)^{\alpha/2}$.] Thus, we seek to find a useful generalization of (5.3) that goes beyond one-parameter stable processes.

We define for all finite Borel measures μ and ν on \mathbf{R}^d , and all kernels κ on \mathbf{R}^d ,

$$(5.4) \quad \mathcal{E}_\kappa(\mu, \nu) := \iint \left(\frac{\kappa(x - y) + \kappa(y - x)}{2} \right) \mu(dx) \nu(dy).$$

This is called the *mutual energy* between μ and ν in gauge κ , and defines a quadratic form with pseudo-norm

$$(5.5) \quad \mathcal{E}_\kappa(\mu) := \mathcal{E}_\kappa(\mu, \mu).$$

This is the “ κ -energy” of the measure μ . There is a corresponding capacity defined as

$$(5.6) \quad \mathcal{C}_\kappa(F) := \frac{1}{\inf \mathcal{E}_\kappa(\mu)},$$

where the infimum is taken over all compactly supported probability measures μ on F , $\inf \emptyset := \infty$, and $1/\infty := 0$.

We can recognize the left-hand side of (5.3) to be $\mathcal{E}_{\kappa_{d-\alpha}}(\mu)$. Because $\kappa_{d-\alpha}(x) < \infty$ if and only if $x \neq 0$, the following is a nontrivial generalization of (5.3).

Theorem 5.2. *Suppose κ is a continuous kernel of positive type on \mathbf{R}^d which satisfies one of the following two conditions:*

- (1) $\kappa(x) < \infty$ if and only if $x \neq 0$;
- (2) $\hat{\kappa} \in L^\infty(\mathbf{R}^d)$, and $\kappa(x) < \infty$ when $x \neq 0$.

Then for all Borel probability measures μ and ν on \mathbf{R}^d ,

$$(5.7) \quad \mathcal{E}_{\kappa*\nu}(\mu) = \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \operatorname{Re} \hat{\nu}(\xi) |\hat{\mu}(\xi)|^2 d\xi.$$

Remark 5.3. If $\kappa \in L^1(\mathbf{R}^d)$, then $\hat{\kappa}$ can be define by the usual Fourier transform, $\hat{\kappa}(\xi) = \int_{\mathbf{R}^d} \exp(ix \cdot \xi) \kappa(x) dx$. That is, the condition $\hat{\kappa} \in L^\infty(\mathbf{R}^d)$ is automatically verified in this case. In fact, $\hat{\kappa}$ is bounded in this case, as can be seen from $\sup_{\xi \in \mathbf{R}^d} |\hat{\kappa}(\xi)| = \|\kappa\|_{L^1(\mathbf{R}^d)}$. \square

Our proof of Theorem 5.2 proceeds in four steps; the first three are stated as lemmas.

The folklore of harmonic analysis contains precise versions of the loose assertion that “typically, kernels of positive type achieve their supremum at the origin.” The first step in the proof of Theorem 5.2 is to verify a suitable form of this statement.

Lemma 5.4. *Suppose κ is a kernel of positive type such that $\hat{\kappa} \in L^\infty(\mathbf{R}^d)$. Suppose also that κ is continuous on all of \mathbf{R}^d , and finite on $\mathbf{R}^d \setminus \{0\}$. Then, $\kappa(0) = \sup_{x \in \mathbf{R}^d} \kappa(x)$.*

Proof. Recall the functions ϕ_ϵ from the proof of Theorem 2.1; see (4.6). Then, $\kappa * \phi_\epsilon \in L^1(\mathbf{R}^d)$ and $\widehat{\kappa * \phi_\epsilon} = \hat{\kappa} \hat{\phi}_\epsilon \geq 0$. But $\operatorname{ess\,sup}_{\xi \in \mathbf{R}^d} |\hat{\kappa}(\xi)| < \infty$ and $\sup_{\xi \in \mathbf{R}^d} |\hat{\phi}_\epsilon(\xi)| \leq \|\phi_\epsilon\|_{L^1(\mathbf{R}^d)} = 1$. Moreover, the construction of ϕ_ϵ ensures that $\hat{\phi}_\epsilon$ is integrable. Therefore, $\widehat{\kappa * \phi_\epsilon} \in L^1(\mathbf{R}^d) \cap L^\infty(\mathbf{R}^d)$, and for Lebesgue-almost all $x \in \mathbf{R}^d$,

$$(5.8) \quad (\kappa * \phi_\epsilon)(x) = \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} e^{-ix \cdot \xi} \hat{\kappa}(\xi) \hat{\phi}_\epsilon(\xi) d\xi,$$

thanks to the inversion formula for Fourier transforms. Since both sides of (5.8) are continuous functions of x , that equation is valid for all $x \in \mathbf{R}^d$. Furthermore, because κ is a kernel of positive type, it follows that $(\kappa * \phi_\epsilon)(x) \leq (2\pi)^{-d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) d\xi$. But if $x \neq 0$, then $\lim_{\epsilon \downarrow 0} (\kappa * \phi_\epsilon)(x) = \kappa(x)$, and hence,

$$(5.9) \quad \sup_{x \in \mathbf{R}^d \setminus \{0\}} \kappa(x) \leq \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) d\xi.$$

It suffices to prove that

$$(5.10) \quad \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) d\xi \leq \kappa(0).$$

This holds trivially if $\kappa(0)$ is infinite. Therefore, we may assume without loss of generality that $\kappa(0) < \infty$. The continuity of κ ensures that it is uniformly continuous in a neighborhood

of the origin, thence we have $\lim_{\epsilon \downarrow 0} (\kappa * \phi_\epsilon)(0) = \kappa(0)$ by the classical Fejér theorem. Also, we recall that $\lim_{\epsilon \downarrow 0} \hat{\phi}_\epsilon(\xi) = 1$ for all $\xi \in \mathbf{R}^d$. We use these facts in conjunction with (5.8) and Fatou's lemma to deduce (5.10), and hence the lemma. \square

Next we present the second step in the proof of Theorem 5.2. This is another folklore fact from harmonic analysis.

We say that a kernel κ is *lower semicontinuous* if there exist a sequence of continuous functions $\kappa_1, \kappa_2, \dots : \mathbf{R}^d \rightarrow \mathbf{R}_+$ such that $\kappa_n(x) \leq \kappa_{n+1}(x)$ for all $n \geq 1$ and $x \in \mathbf{R}^d$, such that $\kappa_n(x) \uparrow \kappa(x)$ for all $x \in \mathbf{R}^d$, as $n \uparrow \infty$. Because $\kappa_n(x)$ is assumed to be in \mathbf{R}_+ [and not $\bar{\mathbf{R}}_+$], our definition of lower semicontinuity is slightly different from the usual one. Nonetheless, the following is a consequence of Lemma 12.11 of Mattila (1995, p. 161).

Lemma 5.5. *Suppose κ is a lower semicontinuous kernel of positive type on \mathbf{R}^d . Then, for all Borel probability measures μ on \mathbf{R}^d ,*

$$(5.11) \quad \mathcal{E}_\kappa(\mu) \leq \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}(\xi)|^2 d\xi.$$

Our next lemma constitutes the third step of our proof of Theorem 5.2.

Lemma 5.6. *Suppose κ is a continuous kernel of positive type on \mathbf{R}^d , which satisfies one of the following two conditions:*

- (1) $\kappa(x) < \infty$ if and only if $x \neq 0$;
- (2) $\hat{\kappa} \in L^\infty(\mathbf{R}^d)$, and $\kappa(x) < \infty$ when $x \neq 0$.

Then, for all Borel probability measures μ on \mathbf{R}^d ,

$$(5.12) \quad \mathcal{E}_\kappa(\mu) = \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}(\xi)|^2 d\xi.$$

A weaker version of this result is stated in Kahane (1985, p. 134) without proof.

According to Kahane (*loc. cit.*), functions κ that satisfy the conditions of Lemma 5.6 are called *potential kernels*. They can be defined, in equivalent terms, as kernels of positive types that are continuous on $\mathbf{R}^d \setminus \{0\}$ and $\lim_{x \rightarrow 0} \kappa(x) = \infty$. We will not use this terminology: the term “potential kernel” is reserved for another object.

Proof of Lemma 5.6. Regardless of whether κ satisfies condition (1) or (2), it is lower semicontinuous. Therefore, in light of Lemma 5.5, it suffices to prove that

$$(5.13) \quad \mathcal{E}_\kappa(\mu) \geq \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}(\xi)|^2 d\xi.$$

From here on, our proof considers two separate cases:

Case 1. First, let us suppose κ satisfies condition (1) of the lemma.

Without loss of generality, we may assume that $\mathcal{E}_\kappa(\mu) < \infty$. Since $\lim_{x \rightarrow 0} \kappa(x) = \infty$, this implies that μ does not charge singletons.

According to the Lusin's theorem, for every $\eta \in (0, 1)$ we can find a compact set $K_\eta \subseteq \mathbf{R}^d \setminus \{0\}$ with $\mu(K_\eta^c) \leq \eta$ such that $\kappa * \mu$ is continuous—and hence uniformly continuous—on K_η . Define for all Borel sets $A \subseteq \mathbf{R}^d$,

$$(5.14) \quad \mu_\eta(A) = \frac{\mu(A \cap K_\eta)}{1 - \eta}.$$

Then μ_η is supported by the compact set K_η , and

$$(5.15) \quad \frac{1}{1 - \eta} \geq \mu_\eta(K_\eta) = \frac{\mu(K_\eta)}{1 - \eta} \geq 1.$$

Now $\lim_{r \rightarrow 0} \mu_\eta(B(x, r)) = 0$ for all $x \in \mathbf{R}^d$, where $B(x, r)$ denotes the ℓ^2 -ball of radius $r > 0$ about $x \in \mathbf{R}^d$. Therefore, a compactness argument reveals that for all $\eta \in (0, 1)$,

$$(5.16) \quad \lim_{r \downarrow 0} \sup_{x \in K_\eta} \mu_\eta(B(x, r)) = 0.$$

For otherwise we can find $\delta > 0$ and $x_r \in K_\eta$ such that for all $r > 0$, $\mu_\eta(B(x_r, r)) \geq \delta$. By compactness we can extract a subsequence $r' \rightarrow 0$ and $x \in K_\eta$ such that $x_{r'} \rightarrow x$. It follows easily then $\mu_\eta(B(x, \epsilon)) \geq \delta$ for all $\epsilon > 0$, whence $\mu_\eta(\{x\}) \geq \delta$, which contradicts the fact that μ_η does not charge singletons.

Next we choose and fix $y \in K_\eta$ and $\eta > 0$. We claim that

$$(5.17) \quad \sup_{x \in K_\eta} \int_{B(x, \epsilon)} \kappa(y - z) \mu_\eta(dz) \rightarrow 0 \quad \text{as } \epsilon \downarrow 0.$$

Indeed, by (5.16) and the fact that $(\kappa * \mu_\eta)(y) < \infty$, we see that for all $\rho > 0$ there exists $\theta > 0$ such that

$$(5.18) \quad \lim_{\epsilon \downarrow 0} \sup_{x \in K_\eta: |x - y| \leq \theta/2} \int_{B(x, \epsilon)} \kappa(y - z) \mu_\eta(dz) \leq \rho.$$

On the other hand, by the continuity of κ on $\mathbf{R}^d \setminus \{0\}$ and (5.16), we have

$$(5.19) \quad \lim_{\epsilon \downarrow 0} \sup_{x \in K_\eta: |x - y| > \theta/2} \int_{B(x, \epsilon)} \kappa(y - z) \mu_\eta(dz) = 0.$$

By combining (5.18) and (5.19), we find that

$$(5.20) \quad \lim_{\epsilon \downarrow 0} \sup_{x \in K_\eta} \int_{B(x, \epsilon)} \kappa(y - z) \mu_\eta(dz) \leq \rho.$$

Thus (5.17) follows from (5.20), because we can choose ρ as small as we want. By (5.17) and another appeal to compactness, we obtain

$$(5.21) \quad \sup_{y \in K_\eta} \sup_{x \in K_\eta} \int_{B(x, \epsilon)} \kappa(y - z) \mu_\eta(dz) \rightarrow 0 \quad \text{as } \epsilon \downarrow 0.$$

Consequently, for all $k \geq 1$, we can find $\epsilon_k \rightarrow 0$ such that

$$(5.22) \quad \sup_{y \in K_\eta} \sup_{x \in K_\eta} \left| (\kappa * \mu_\eta)(y) - \int_{B(x, \epsilon_k)^c} \kappa(y - z) \mu_\eta(dz) \right| \leq \frac{1}{k}.$$

Let $y \in K_\eta$ and let $\{y_n\}_{n=1}^\infty$ be an arbitrary sequence in K_η such that $\lim_{n \rightarrow \infty} y_n = y$. Because $z \mapsto \kappa(y - z)$ is uniformly continuous on $B(y, \epsilon_k)^c \cap K_\eta$, we have

$$(5.23) \quad \lim_{n \rightarrow \infty} \int_{B(y, \epsilon_k)^c} \kappa(y_n - z) \mu_\eta(dz) = \int_{B(y, \epsilon_k)^c} \kappa(y - z) \mu_\eta(dz).$$

This and (5.22) together imply that for all $k \geq 1$,

$$(5.24) \quad \lim_{n \rightarrow \infty} |(\kappa * \mu_\eta)(y_n) - (\kappa * \mu_\eta)(y)| \leq \frac{2}{k}.$$

Let $k \uparrow \infty$ to deduce that $\kappa * \mu_\eta$ is continuous, and hence uniformly continuous, on K_η . On the other hand, it can be verified directly that $\kappa * \mu_\eta$ is continuous on K_η^c . Hence, we have shown that $\kappa * \mu_\eta$ is continuous on \mathbf{R}^d .

If $0 < \epsilon, \eta < 1$, then we can appeal to the Fubini-Tonelli theorem, a few times in succession, to deduce that

$$(5.25) \quad \mathcal{E}_{\kappa * \psi_\epsilon}(\mu_\eta) = \int \kappa(-z) (a_{\epsilon, \eta} * b_{\epsilon, \eta})(z) dz,$$

where $a_{\epsilon, \eta} := \phi_\epsilon * \mu_\eta$, $b_{\epsilon, \eta} := \phi_\epsilon * \check{\mu}_\eta$, and $\check{\mu}_\eta$ is the Borel probability measure on \mathbf{R}^d that is defined by

$$(5.26) \quad \check{\mu}_\eta(A) := \mu_\eta(-A) \quad \text{for all } A \subseteq \mathbf{R}^d.$$

We observe the following elementary facts:

- (1) Both $a_{\epsilon, \eta}$ and $b_{\epsilon, \eta}$ are infinitely differentiable functions of compact support;
- (2) because ϕ_ϵ is of positive type, the Fourier transform of $a_{\epsilon, \eta} * b_{\epsilon, \eta}$ is $|\hat{\phi}_\epsilon|^2 |\hat{\mu}_\eta|^2$;
- (3) the Fourier transform of $z \mapsto \kappa(-z)$ is the same as that of κ because κ is of positive type;

Therefore, we can combine the preceding with the Parseval identity, and deduce that

$$(5.27) \quad \begin{aligned} \mathcal{E}_{\kappa * \psi_\epsilon}(\mu_\eta) &= \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \overline{\hat{a}_{\epsilon, \eta}(\xi)} \hat{b}_{\epsilon, \eta}(\xi) d\xi \\ &= \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\phi}_\epsilon(\xi)|^2 |\hat{\mu}_\eta(\xi)|^2 d\xi. \end{aligned}$$

We can apply the Fubini-Tonelli theorem to write the left-most term in another way, as well. Namely,

$$(5.28) \quad \mathcal{E}_{\kappa * \psi_\epsilon}(\mu_\eta) = \int (\kappa * \psi_\epsilon * \mu_\eta) d\mu_\eta.$$

The continuity of $\kappa * \mu_\eta$, and Fejér's theorem, together imply that $\kappa * \mu_\eta * \psi_\epsilon$ converges to $\kappa * \mu_\eta$ uniformly on K_η as $\epsilon \downarrow 0$, and hence $\mathcal{E}_{\kappa * \psi_\epsilon}(\mu_\eta)$ converges to $\mathcal{E}_\kappa(\mu_\eta)$ as $\epsilon \downarrow 0$. This and (5.27) together imply that

$$(5.29) \quad \begin{aligned} \mathcal{E}_\kappa(\mu_\eta) &= \frac{1}{(2\pi)^d} \lim_{\epsilon \rightarrow 0} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\phi}_\epsilon(\xi)|^2 |\hat{\mu}_\eta(\xi)|^2 d\xi \\ &\geq \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}_\eta(\xi)|^2 d\xi, \end{aligned}$$

owing to Fatou's lemma. Since μ_η is $(1 - \eta)^{-1}$ times a restriction of μ , it follows that

$$(5.30) \quad \frac{1}{(1 - \eta)^2} \mathcal{E}_\kappa(\mu) \geq \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}_\eta(\xi)|^2 d\xi.$$

Because for all $\xi \in \mathbf{R}^d$,

$$(5.31) \quad \left| \hat{\mu}_\eta(\xi) - \frac{\hat{\mu}(\xi)}{1 - \eta} \right| \leq \frac{\mu(K_\eta^c)}{1 - \eta} \leq \frac{\eta}{1 - \eta},$$

we deduce that $\lim_{\eta \rightarrow 0} \hat{\mu}_\eta = \hat{\mu}$ pointwise. An appeal to Fatou's lemma justifies (5.13), and this completes our proof in the first case.

Case 2. Next, we consider the case that κ satisfies condition (2) of the lemma. If $\kappa(0) = \infty$, then condition (1) is satisfied, and therefore the proof is complete. Thus, we may assume that $\kappa(0) < \infty$. According to Lemma 5.4, κ is a bounded and continuous function from \mathbf{R}^d into \mathbf{R}_+ . For all Borel sets $A \subseteq \mathbf{R}^d$ define

$$(5.32) \quad \mu_n(A) := \frac{\mu_n(A \cap [-n, n]^d)}{\chi_n}, \quad \text{where} \quad \chi_n := \mu([-n, n]^d).$$

If $n > 0$ is sufficiently large, then μ_n is a well-defined Borel probability measure on $[-n, n]^d$, and since κ is uniformly continuous on $[-n, n]^d$,

$$(5.33) \quad \begin{aligned} \frac{1}{\chi_n^2} \mathcal{E}_\kappa(\mu) &\geq \mathcal{E}_\kappa(\mu_n) \\ &= \lim_{\epsilon \downarrow 0} \mathcal{E}_{\kappa * \psi_\epsilon}(\mu_n) \\ &= \frac{1}{(2\pi)^d} \lim_{\epsilon \downarrow 0} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \hat{\psi}_\epsilon(\xi) |\hat{\mu}_n(\xi)|^2 d\xi. \end{aligned}$$

The first equality holds because $\kappa * \psi_\epsilon$ converges uniformly to κ on $[-n, n]^d$, as $\epsilon \downarrow 0$. The second is a consequence of the Parseval identity. Thanks to Fatou's lemma, we have proved that for all $n > 0$ sufficiently large,

$$(5.34) \quad \frac{1}{\chi_n^2} \mathcal{E}_\kappa(\mu) \geq \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}_n(\xi)|^2 d\xi.$$

As n tends to infinity, $\chi_n \uparrow 1$ and $\hat{\mu}_n \rightarrow \hat{\mu}$ pointwise. Therefore, another appeal to Fatou's lemma implies (5.13), and hence the lemma. \square

Now we derive Theorem 5.2.

Proof of Theorem 5.2. Let us observe that by the Fubini-Tonelli theorem,

$$(5.35) \quad \mathcal{E}_{\kappa*\nu}(\mu) = \mathcal{E}_\kappa(\mu, \check{\nu} * \mu).$$

In fact, both sides are equal to $E[\kappa(X - X' - Y)]$, where (X, X', Y) are independent, X and X' are distributed as μ , and Y is distributed as ν .

Lemma 5.6, and polarization, together imply that the following holds for every Borel probability measure σ on \mathbf{R}^d :

$$(5.36) \quad \mathcal{E}_\kappa(\mu, \sigma) = \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \hat{\sigma}(\xi) \overline{\hat{\mu}(\xi)} d\xi.$$

Indeed, we first notice that

$$(5.37) \quad \mathcal{E}_\kappa\left(\frac{\mu + \sigma}{2}\right) = \frac{1}{4} \mathcal{E}_\kappa(\mu) + \frac{1}{4} \mathcal{E}_\kappa(\sigma) + \frac{1}{2} \mathcal{E}_\kappa(\mu, \sigma).$$

Thus, we solve for $\mathcal{E}_\kappa(\mu, \sigma)$ and apply Lemma 5.6 to deduce that

$$(5.38) \quad \begin{aligned} \mathcal{E}_\kappa(\mu, \sigma) &= \frac{2}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \left| \widehat{\left(\frac{\mu + \sigma}{2}\right)}(\xi) \right|^2 d\xi \\ &\quad - \frac{1}{2(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\mu}(\xi)|^2 d\xi - \frac{1}{2(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) |\hat{\sigma}(\xi)|^2 d\xi. \end{aligned}$$

We solve to obtain (5.36).

Thus we define $\sigma := \check{\nu} * \mu$, observe that σ is a Borel probability measure on \mathbf{R}^d , and $\mathcal{E}_{\kappa*\nu}(\mu) = \mathcal{E}_\kappa(\mu, \sigma)$. Thus, we may apply (5.35) and (5.36)—in this order—to find that

$$(5.39) \quad \begin{aligned} \mathcal{E}_{\kappa*\nu}(\mu) &= \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \hat{\sigma}(\xi) \overline{\hat{\mu}(\xi)} d\xi \\ &= \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \hat{\kappa}(\xi) \hat{\nu}(-\xi) |\hat{\mu}(\xi)|^2 d\xi. \end{aligned}$$

In the last line, we have only used the fact that $\hat{\sigma}(\xi) = \hat{\nu}(-\xi)\hat{\mu}(\xi)$ for all $\xi \in \mathbf{R}^d$. Because the left-most term in (5.39) is real-valued, so is the right-most term. Therefore, we may consider only the real part of the right-most item in (5.39), and this proves the result. \square

6. Absolute continuity considerations

Suppose E is a Borel measurable subset of \mathbf{R}^N that has positive Lebesgue measure, and $Y := \{Y(\mathbf{t})\}_{\mathbf{t} \in E}$ is an \mathbf{R}^d -valued random field that is indexed by E .

Definition 6.1. We say that Y has a *one-potential density* v if $v \in L^1(\mathbf{R}^d)$ is nonnegative and satisfies the following for all Borel measurable functions $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$:

$$(6.1) \quad \frac{1}{\int_E e^{-[s]} d\mathbf{s}} \mathbb{E} \left[\int_E f(Y(\mathbf{t})) e^{-[t]} d\mathbf{t} \right] = \int_{\mathbf{R}^d} f(x) v(x) dx.$$

In particular, $\tilde{\mathfrak{X}}$ has a one-potential density v if for all Borel measurable functions $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$,

$$(6.2) \quad (Rf)(x) = \int_{\mathbf{R}^d} f(x+y) v(y) dy := (f * \check{v})(x) \quad \text{for all } x \in \mathbf{R}^d,$$

where $\check{g}(z) := g(-z)$ for all functions g . It follows from Lemma 2.1 in Hawkes (1979) that $\tilde{\mathfrak{X}}$ has a one-potential density if and only if the operator R defined by (2.9) is *strong Feller*. That is, if Rf is continuous whenever f is Borel measurable and has compact support.

In order to be concrete, we choose a “nice” version of the one-potential density v , when it exists. Before we proceed further, let us observe that when v exists it is a probability density. In particular, the Lebesgue density theorem tell us that

$$(6.3) \quad v(x) = \lim_{\epsilon \downarrow 0} \frac{1}{\lambda_d(B(0, \epsilon))} \int_{B(x, \epsilon)} v(y) dy \quad \text{for almost all } x \in \mathbf{R}^d [\lambda_d].$$

We can recognize the integral as $(R\mathbf{1}_{B(0, \epsilon)})(-x)$. Moreover, we can alter v on a Lebesgue-null set and still obtain a one-potential density for $\tilde{\mathfrak{X}}$. Therefore, from now on, we *always* choose the following version of v :

$$(6.4) \quad v(x) := \liminf_{\epsilon \downarrow 0} \frac{(R\mathbf{1}_{(0, \epsilon)})(-x)}{\lambda_d(B(0, \epsilon))} \quad \text{for all } x \in \mathbf{R}^d.$$

Lemma 2.3 states that the Fourier multiplier of R is K_{Ψ} . Therefore, the Fourier transform of \check{v} is also K_{Ψ} ; confer with (6.2). Because K_{Ψ} is real-valued, this proves the following:

Lemma 6.2. *If $\tilde{\mathfrak{X}}$ has a one-potential density v , then $\hat{v} = K_{\Psi}$ in the sense of Schwartz. In particular, v is an integrable kernel of positive type on \mathbf{R}^d .*

Remark 6.3. Let $\epsilon_1, \dots, \epsilon_N$ denote N random variables, all independent of one another, as well as $\{X_j\}_{j=1}^N$ and $\{X'_j\}_{j=1}^N$, with $\mathbb{P}\{\epsilon_1 = \pm 1\} = 1/2$. We can re-organize the order of

integration a few times to find that

$$(6.5) \quad \frac{1}{2^N} \mathbb{E} \left[\int_{\mathbf{R}^N} f(\tilde{\mathfrak{X}}(\mathbf{t})) e^{-[\mathbf{t}]} d\mathbf{t} \right] = \mathbb{E} \left[\int_{\mathbf{R}_+^N} f \left(\sum_{j=1}^N \epsilon_j X_j(s_j) \right) e^{-[\mathbf{s}]} d\mathbf{s} \right].$$

Thus, in particular, $\tilde{\mathfrak{X}}$ has a one-potential density if and only if the \mathbf{R}_+^N -indexed random field $(t_1, \dots, t_N) \mapsto \sum_{j=1}^N \epsilon_j X_j(t_j)$ does [interpreted in the obvious sense]. Interestingly enough, the latter random field appears earlier—though for quite different reasons as ours—in the works of Marcus and Rosen (1999b, 1999a). \square

Remark 6.4. It follows from the definitions that a random field $\{Y(\mathbf{t})\}_{\mathbf{t} \in E}$ has a one-potential density if and only if the Borel probability measure

$$(6.6) \quad \mathbf{R}^d \supseteq A \mapsto \mathbb{E} \left[\int_E \mathbf{1}_A(Y(\mathbf{t})) e^{-[\mathbf{t}]} d\mathbf{t} \right] \quad \text{is absolutely continuous.}$$

Now suppose $f \geq 0$ is Borel measurable. Then

$$(6.7) \quad \int_{\mathbf{R}^N} f(\tilde{\mathfrak{X}}(\mathbf{t})) e^{-[\mathbf{t}]} d\mathbf{t} \geq \int_{\mathbf{R}_+^N} f(\mathfrak{X}(\mathbf{t})) e^{-[\mathbf{t}]} d\mathbf{t}.$$

We take expectations of both sides to deduce from (6.6) that if $\tilde{\mathfrak{X}}$ has a one-potential density, then so does \mathfrak{X} . When $N = 1$, it can be verified that the converse is also true. However, the previous remark can be used to show that, when $N \geq 2$, the converse is not necessarily true. Thus, we can conclude that the existence of a one-potential density for $\tilde{\mathfrak{X}}$ is a *more stringent* condition than the existence of a one-potential density for \mathfrak{X} . Another consequence of the preceding is the following: If the one-potential density of $\tilde{\mathfrak{X}}$ exists and the potential density of \mathfrak{X} is a.e.-positive, then the one-potential density of $\tilde{\mathfrak{X}}$ is per force also a.e.-positive. We have, and will, encounter these conditions several times. \square

Our next theorem is the main result of this subsection. From a technical point of view, it is also a key result in this paper. In order to describe it properly, we introduce some notation first.

If T is a nonempty subset of $\{1, \dots, N\}$, then we define $|T|$ to be the cardinality of T , and \mathfrak{X}_T to be the *subprocess* associated to the index T . That is, \mathfrak{X}_T is the following $|T|$ -parameter, \mathbf{R}^d -valued random field:

$$(6.8) \quad \mathfrak{X}_T(\mathbf{t}) := \sum_{j \in T} X_j(t_j) \quad \text{for all } \mathbf{t} \in \mathbf{R}_+^{|T|}.$$

In order to obtain nice formulas, we define \mathfrak{X}_\emptyset to be the constant 0. In this way, it follows that, regardless of whether or not $|T| > 0$, each \mathfrak{X}_T is itself an additive Lévy process, and the Lévy exponent of \mathfrak{X}_T is the function $(\Psi_j)_{j \in T} : \mathbf{R}^d \rightarrow \mathbf{C}^{|T|}$. Thus, we can talk about the

stationary field $\tilde{\mathfrak{X}}_T$ for all $T \subseteq \{1, \dots, N\}$, etc. Despite this new notation, we continue to write \mathfrak{X} and $\tilde{\mathfrak{X}}$ in place of the more cumbersome $\mathfrak{X}_{\{1, \dots, N\}}$ and $\tilde{\mathfrak{X}}_{\{1, \dots, N\}}$.

Theorem 6.5. *Suppose there exists a nonrandom and nonempty subset $T \subseteq \{1, \dots, N\}$, such that $\tilde{\mathfrak{X}}_T$ has a one-potential density $v_T : \mathbf{R}^d \mapsto \bar{\mathbf{R}}_+$ that is continuous on \mathbf{R}^d , and finite on $\mathbf{R}^d \setminus \{0\}$. Then, $\tilde{\mathfrak{X}}$ has a one-potential density v on \mathbf{R}^d , and for all Borel sets $F \subseteq \mathbf{R}^d$,*

$$(6.9) \quad \mathbb{E} [\lambda_d (\mathfrak{X}(\mathbf{R}_+^N) \oplus F)] > 0 \quad \Leftrightarrow \quad \mathcal{C}_v(F) > 0.$$

Proof. First, consider the case that $T = \{1, \dots, N\}$. In this case, $\tilde{\mathfrak{X}}$ has a one-potential density $v := v_{\{1, \dots, N\}}$ that is continuous on \mathbf{R}^d , and finite away from the origin. In light of Remark 5.3, Lemma 6.2 shows us that $\kappa := v$ satisfies condition (2) of Theorem 5.2. Theorem 5.2, in turn, implies that $I_\Psi(\mu) = \mathcal{C}_v(\mu)$ for all Borel probability measures μ on \mathbf{R}^d , and thence $\mathcal{C}_v(F) = \text{cap}_\Psi(F)$. This and Theorem 1.1 together imply Theorem 6.5 in the case that $\tilde{\mathfrak{X}}$ has a continuous one-potential density that is finite on $\mathbf{R}^d \setminus \{0\}$.

Next, consider the remaining case that $1 \leq |T| \leq N - 1$. For all Borel sets $A \subseteq \mathbf{R}^d$ define

$$(6.10) \quad V_{T^c}(A) := \mathbb{E} \left[\int_{\mathbf{R}^{N-|T|}} \mathbf{1}_A \left(\tilde{\mathfrak{X}}_{T^c}(\mathbf{s}) \right) e^{-[\mathbf{s}]} d\mathbf{s} \right].$$

This is the one-potential measure for the stationary field based on the additive Lévy process $\tilde{\mathfrak{X}}_{T^c}$. The corresponding object for $\tilde{\mathfrak{X}}_T$ can be defined likewise, viz.,

$$(6.11) \quad V_T(A) := \mathbb{E} \left[\int_{\mathbf{R}^{|T|}} \mathbf{1}_A \left(\tilde{\mathfrak{X}}_T(\mathbf{t}) \right) e^{-[\mathbf{t}]} d\mathbf{t} \right].$$

We can observe that

$$(6.12) \quad \begin{aligned} V_{T^c}(A) &:= \mathbb{E} \left[\int_{\mathbf{R}_+^N} \mathbf{1}_A \left(\sum_{j \in T^c} \epsilon_j X_j(s_j) \right) e^{-[\mathbf{s}]} d\mathbf{s} \right], \\ V_T(A) &:= \mathbb{E} \left[\int_{\mathbf{R}_+^N} \mathbf{1}_A \left(\sum_{j \in T} \epsilon_j X_j(s_j) \right) e^{-[\mathbf{s}]} d\mathbf{s} \right], \end{aligned}$$

where $\epsilon_1, \dots, \epsilon_N$ are i.i.d. random functions that are totally independent of X_1, \dots, X_N and $\tilde{X}_1, \dots, \tilde{X}_N$, and $\mathbb{P}\{\epsilon_1 = \pm 1\} = 1/2$.

Similarly, we can write for all Borel measurable functions $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$ and $x \in \mathbf{R}^d$,

$$(6.13) \quad (Rf)(x) = \mathbb{E} \left[\int_{\mathbf{R}_+^N} f \left(x + \sum_{j=1}^N \epsilon_j X_j(t_j) \right) e^{-[\mathbf{t}]} d\mathbf{t} \right].$$

Note that the integral is now over \mathbf{R}_+^N [and not \mathbf{R}^N]. Moreover, we can write

$$(6.14) \quad \begin{aligned} (Rf)(x) &= \iint f(x+y+z) V_T(dy) V_{T^c}(dz) \\ &= \int f(x+y) (V_T * V_{T^c})(dx), \end{aligned}$$

The condition that $\tilde{\mathfrak{X}}_T$ has a one-potential density v_T is equivalent to the statement that $V_T(dy) = v_T(y) dy$, whence it follows that $(Rf)(x) = \int f(x+y)v(y) dy$, where

$$(6.15) \quad v(y) := \int v_T(y-z)V_{T^c}(dz).$$

This proves the assertion that $\tilde{\mathfrak{X}}$ has a one-potential density v . Note that v is of the form $\kappa * \nu$, where $\kappa := v_T$ and $\nu := V_{T^c}$. Thus, another appeal to Theorem 5.2 shows that condition (2) there is satisfied [confer also with Remark 5.3], and thence it follows that $\mathcal{C}_v(F) = \text{cap}_{\Psi}(F)$. This and Theorem 1.1 together complete the proof. \square

6.1. A relation to an intersection problem. Define for all Borel sets $F \subseteq \mathbf{R}^d$,

$$(6.16) \quad \begin{aligned} H_F &:= \{x \in \mathbf{R}^d : \mathbf{P} \{ \mathfrak{X}((0, \infty)^N) \cap (x \oplus F) \neq \emptyset \} > 0 \} \\ \tilde{H}_F &:= \{x \in \mathbf{R}^d : \mathbf{P} \{ \tilde{\mathfrak{X}}(\mathbf{R}_{\neq}^N) \cap (x \oplus F) \neq \emptyset \} > 0 \}, \end{aligned}$$

where $\mathbf{R}_{\neq}^N := \{t \in \mathbf{R}^N : t_j \neq 0 \text{ for every } j = 1, \dots, N\}$. Clearly, $H_F \subseteq \tilde{H}_F$.

The following improves on Proposition 6.2 of Khoshnevisan, Xiao, and Zhong (2003).

Proposition 6.6. *Consider the following statements:*

- (0) $\mathbf{E}[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N) \oplus F)] > 0$;
- (1) $\text{cap}_{\Psi}(F) > 0$;
- (2) $\mathcal{C}_v(F) > 0$;
- (3) $\lambda_d(H_F) > 0$;
- (4) $H_F \neq \emptyset$;
- (5) $\tilde{H}_F = \mathbf{R}^d$; and
- (6) $H_F = \mathbf{R}^d$.

Then:

- (a) *It is always the case that (0) \Leftrightarrow (1).*
- (b) *If \mathfrak{X} has a one-potential density, then (1) \Leftrightarrow (3) \Leftrightarrow (4).*
- (c) *If \mathfrak{X} has an a.e.-positive one-potential density, then (1) \Leftrightarrow (3) \Leftrightarrow (4) \Leftrightarrow (6).*
- (d) *If there exists a non-empty $T \subseteq \{1, \dots, N\}$ such that the sub-process $\tilde{\mathfrak{X}}_T$ has a one-potential density v_T , and v_T is continuous on \mathbf{R}^d and finite on $\mathbf{R}^d \setminus \{0\}$. Then, (1) \Leftrightarrow (2) \Leftrightarrow (3) \Leftrightarrow (4).*

(e) Suppose the conditions of (e) are met, and let v denote the one-potential density of $\tilde{\mathfrak{X}}$. If $v > 0$ a.e., then $(1) \Leftrightarrow (2) \Leftrightarrow (3) \Leftrightarrow (4) \Leftrightarrow (5)$.

Remark 6.7. Proposition 6.6 is deceptively subtle. For example, The preceding proposition might fail to hold if H_F were replaced by the related set

$$(6.17) \quad H_F^* := \{x \in \mathbf{R}^d : \mathbf{P} \{ \mathfrak{X}(\mathbf{R}_+^N) \cap (x \oplus F) \neq \emptyset \} > 0 \}.$$

Indeed, because $\mathfrak{X}(\mathbf{0}) = 0$, it follows that $H_F^* = -F$. \square

Now we prove Proposition 6.6.

Proof of Proposition 6.6. Part (a) is merely a restatement of Theorem 1.1.

We first observe the following consequence of the Fubini-Tonelli theorem:

$$(6.18) \quad \mathbf{E} [\lambda_d(\mathfrak{X}((0, \infty)^N) \ominus F)] = \int_{\mathbf{R}^d} \mathbf{P} \{ \mathfrak{X}((0, \infty)^N) \cap (x \oplus F) \neq \emptyset \} dx.$$

See (2.7) for a similar computation. If $\text{cap}_{\Psi}(F) > 0$, then the left-hand side of (6.18) is positive by a simple variant of Theorem 1.1. This proves that $(1) \Rightarrow (3)$. Conversely, if $\lambda_d(H_F) > 0$, then the right-hand side of (6.18)—and hence the left-hand side—are positive. Another appeal to Theorem 1.1 shows us that $(1) \Leftrightarrow (3)$.

Now we finish the proof of (b). Clearly, (3) implies (4). In order to prove that $(4) \Rightarrow (3)$, we define

$$(6.19) \quad Q_{\mathbf{s}} := (s_1, \infty) \times \cdots \times (s_N, \infty) \quad \text{for all } \mathbf{s} \in \mathbf{R}^N.$$

We note that for all $\mathbf{s} \in \mathbf{R}^N$ and $x \in \mathbf{R}^d$,

$$(6.20) \quad \begin{aligned} & \mathbf{P} \{ \tilde{\mathfrak{X}}(Q_{\mathbf{s}}) \cap (x \oplus F) \neq \emptyset \} \\ &= \int_{\mathbf{R}^d} \mathbf{P} \{ \mathfrak{X}((0, \infty)^N) \cap ((x - y) \oplus F) \neq \emptyset \} \mathbf{P} \{ \tilde{\mathfrak{X}}(\mathbf{s}) \in dy \}. \end{aligned}$$

This uses only the fact that $\{\tilde{\mathfrak{X}}(\mathbf{t}) - \tilde{\mathfrak{X}}(\mathbf{s})\}_{\mathbf{t} \in Q_{\mathbf{s}}}$ is independent of $\tilde{\mathfrak{X}}(\mathbf{s})$, and has the same law as \mathfrak{X} . In particular,

$$(6.21) \quad \begin{aligned} & \int_{\mathbf{R}^N} \mathbf{P} \{ \tilde{\mathfrak{X}}(Q_{\mathbf{s}}) \cap (x \oplus F) \neq \emptyset \} e^{-[\mathbf{s}]} d\mathbf{s} \\ &= 2^N \int_{\mathbf{R}^d} \mathbf{P} \{ \mathfrak{X}((0, \infty)^N) \cap ((x - y) \oplus F) \neq \emptyset \} v(y) dy, \end{aligned}$$

where v denotes the one-potential density of $\tilde{\mathfrak{X}}$. [It exists thanks to Remark 6.4.]

If (4) holds, then there exists $\mathbf{s} \in \mathbf{R}_+^N$ such that $\mathbf{P} \{ \mathfrak{X}(Q_{\mathbf{s}}) \cap (x \oplus F) \neq \emptyset \} > 0$. Consequently, $\mathbf{P} \{ \tilde{\mathfrak{X}}(Q_{\mathbf{s}}) \cap (x \oplus F) \neq \emptyset \} > 0$. As we decrease \mathbf{s} , coordinatewise, the preceding probability increases. Therefore, the left-most term in (6.21) is (strictly) positive, and therefore

so is the right-most term in (6.21). We can conclude that $P\{\mathfrak{X}((0, \infty)^N) \cap (w \oplus F) \neq \emptyset\} > 0$ for all w in a non-null set $[\lambda_d]$. This proves (3) \Rightarrow (4), whence (b).

In order to prove (c), let g denote the one-potential density of \mathfrak{X} . Then, we argue as we did to prove (6.21), and deduce that

$$(6.22) \quad \int_{\mathbf{R}_+^N} P\{\mathfrak{X}(Q_s) \cap (x \oplus F) \neq \emptyset\} e^{-[s]} d\mathbf{s} \\ = \int_{\mathbf{R}^d} P\{\mathfrak{X}((0, \infty)^N) \cap ((x - y) \oplus F) \neq \emptyset\} g(y) dy.$$

If $H_F \neq \mathbf{R}^d$, then the left-most term must be zero for some $x \in \mathbf{R}^d$. Because $g > 0$ a.e., this proves that $\lambda_d(H_F) = 0$.

The equivalence of (1) \Leftrightarrow (2) in (d) is contained in Theorems 1.1 and 6.5. The rest of Part (d) is proved similarly as in the proof of (b).

To prove (e): We note that “ $\widetilde{H}_F = \mathbf{R}^d$ iff (4)” [under the condition that $v > 0$ a.e.] has a very similar proof to “ $H_F = \mathbf{R}^d$ iff (4)” [under the condition that $g > 0$ a.e.], but uses (6.21) instead of (6.22). We omit the details and conclude our proof. \square

7. Proof of Theorem 1.3

We continue to let X_1, \dots, X_N denote N independent Lévy processes on \mathbf{R}^d with respective exponents Ψ_1, \dots, Ψ_N . There is a large literature that is devoted to the “ N -parameter Lévy process” $\otimes_{j=1}^N X_j$ defined by

$$(7.1) \quad (\otimes_{j=1}^N X_j)(\mathbf{t}) := \begin{pmatrix} X_1(t_1) \\ \vdots \\ X_N(t_N) \end{pmatrix} \quad \text{for all } \mathbf{t} \in \mathbf{R}_+^N.$$

Note that:

- (i) The state space of $\otimes_{j=1}^N X_j$ is $(\mathbf{R}^d)^N$; and
- (ii) In the special case that $N = 2$, (7.1) reduces to (1.2).

In this section we describe how this theory, and much more, is contained within the theory of additive Lévy processes of this paper.

Let us begin by making the observation that product Lévy processes are in fact degenerate additive Lévy processes. Indeed, for all $\mathbf{t} \in \mathbf{R}_+^N$ we can write

$$(7.2) \quad (\otimes_{j=1}^N X_j)(\mathbf{t}) = \mathbf{A}^1 X_1(t_1) + \dots + \mathbf{A}^N X_N(t_N),$$

where each $X_j(t_j)$ is viewed as a column vector with dimension d , and each \mathbf{A}^j is the $Nd \times d$ matrix

$$(7.3) \quad \mathbf{A}^j = \begin{pmatrix} \mathbf{0}_{(j-1)d \times d} \\ \mathbf{I}_{d \times d} \\ \mathbf{0}_{(N-j)d \times d} \end{pmatrix}.$$

Here: (i) $\mathbf{0}_{(j-1)d \times d}$ is a $(j-1)d \times d$ matrix of zeros; and (ii) $\mathbf{I}_{d \times d}$ is the identity matrix with d^2 entries. We emphasize that the matrix $\mathbf{I}_{d \times d}$ appears after $j-1$ square matrices of size d^2 whose entries are all zeroes.

It is also easy to describe the law of the process $\otimes_{j=1}^N X_j$ via the following characteristic-function relation:

$$(7.4) \quad \mathbb{E} \left[\exp \left\{ i \sum_{j=1}^N \xi^j \cdot X_j(t_j) \right\} \right] = \exp \left(- \sum_{j=1}^N t_j \Psi_j(\xi^j) \right),$$

valid for all $\xi^1, \dots, \xi^N \in \mathbf{R}^d$ and $t_1, \dots, t_N \geq 0$. Therefore, the exponent of the N -parameter additive Lévy process $\otimes_{j=1}^N X_j$ is

$$(7.5) \quad \Psi(\xi^1, \dots, \xi^N) = \begin{pmatrix} \Psi_1(\xi^1) \\ \vdots \\ \Psi_N(\xi^N) \end{pmatrix} \quad \text{for all } \xi^1, \dots, \xi^N \in \mathbf{R}^d.$$

Therefore, Theorem 1.1 applies, without further thought, to yield the following.

Corollary 7.1. *Choose and fix a Borel set $F \subseteq (\mathbf{R}^d)^N$. Then $(\otimes_{j=1}^N X_j)(\mathbf{R}_+^N) \oplus F$ has positive Lebesgue measure with positive probability if and only if there exists a compact-support Borel probability measure μ on F such that*

$$(7.6) \quad \int_{(\mathbf{R}^d)^N} |\hat{\mu}(\xi^1, \dots, \xi^N)|^2 \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi^j)} \right) d\xi < \infty.$$

Next let us suppose that X_1, \dots, X_N have one-potential densities u_1, \dots, u_N , respectively. According to Definition 6.1,

$$(7.7) \quad \mathbb{E} \left[\int_0^\infty f(X_j(s)) e^{-s} ds \right] = \int_{\mathbf{R}^d} f(a) u_j(a) da,$$

for all $j = 1, \dots, N$, and all Borel measurable functions $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$. [This agrees with the usual nomenclature of probabilistic potential theory.]

Lemma 7.2. *If X_1, \dots, X_N have respective one-potential densities u_1, \dots, u_N , then $\otimes_{j=1}^N X_j$ has the one-potential density $u(x^1, \dots, x^N) := \prod_{j=1}^N u_j(x^j)$ for all $x^1, \dots, x^N \in \mathbf{R}^d$. Also,*

$\widetilde{\otimes_{j=1}^N X_j}$ has the one-potential density

$$(7.8) \quad v(x^1, \dots, x^N) := \prod_{j=1}^N \left(\frac{u_j(x^j) + u_j(-x^j)}{2} \right) \quad \text{for all } x^1, \dots, x^N \in \mathbf{R}^d.$$

Finally, if $u_j(0) > 0$ for all j , then u and v are strictly positive everywhere.

Proof. Let us prove the first assertion about the one-potential density of $\otimes_{j=1}^N X_j$ only. The second assertion is proved very similarly.

We seek to establish that for all Borel measurable functions $f : (\mathbf{R}^d)^N \rightarrow \mathbf{R}_+$,

$$(7.9) \quad \mathbf{E} \left[\int_{\mathbf{R}_+^N} f((\otimes_{j=1}^N X_j)(t)) e^{-[t]} dt \right] = \int_{(\mathbf{R}^d)^N} f(x^1, \dots, x^N) \prod_{j=1}^N u_j(x^j) dx.$$

A density argument reduces the problem to the case that $f(x^1, \dots, x^N)$ has the special form $\prod_{j=1}^N f_j(x^j)$, where $f_j : \mathbf{R}^d \rightarrow \mathbf{R}_+$ is Borel measurable. But in this case, the claim follows immediately from (7.7) and the independence of X_1, \dots, X_N .

In order to complete the proof, consider the case that $u_j(0) > 0$. Lemma 3.2 of Evans (1987b) asserts that $u_j(z) > 0$ for all $z \in \mathbf{R}^d$, whence follows the lemma. \square

The preceding lemma and Proposition 6.6 together imply, without any further effort, the following two theorems.

Corollary 7.3. *Let X_1, \dots, X_N be independent Lévy processes on \mathbf{R}^d , and assume that each X_j has a one-potential density u_j such that $u_j(0) > 0$. Then, for all Borel sets $F \subseteq (\mathbf{R}^d)^N$,*

$$(7.10) \quad \mathbf{P} \{ (\otimes_{j=1}^N X_j) ((0, \infty)^N) \cap F \neq \emptyset \} > 0$$

if and only if there exists a compact-support Borel probability measure μ on F such that

$$(7.11) \quad \int_{(\mathbf{R}^d)^N} |\hat{\mu}(\xi^1, \dots, \xi^N)|^2 \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi^j)} \right) d\xi < \infty.$$

Corollary 7.4 (Fitzsimmons and Salisbury, 1989). *Suppose, in addition to the hypotheses of Corollary 7.3, that each $u_j : \mathbf{R}^d \rightarrow \bar{\mathbf{R}}_+$ is continuous on \mathbf{R}^d , and finite on $\mathbf{R}^d \setminus \{0\}$. Then, for all Borel sets $F \subseteq (\mathbf{R}^d)^N$, (7.10) holds if and only if there exists a compact-support Borel probability measure μ on F such that*

$$(7.12) \quad \iint \prod_{j=1}^N \left(\frac{u_j(x^j - y^j) + u_j(y^j - x^j)}{2} \right) \mu(dx^1 \dots dx^N) \mu(dy^1 \dots dy^N) < \infty.$$

We can now prove Theorem 1.3 of the Introduction.

Proof of Theorem 1.3. We observe that:

- (i) $\cap_{j=1}^N X_j((0, \infty))$ can intersect a set $F \subseteq \mathbf{R}^d$ if and only if $(\otimes_{j=1}^N X_j)((0, \infty)^N)$ can intersect $G := \{x \otimes \cdots \otimes x : x \in F\} \subseteq (\mathbf{R}^d)^N$; and
- (ii) any compact-support Borel probability measure σ on G manifestly has the form $\sigma(dx^1 \cdots dx^N) = \mu(dx^1) \prod_{j=2}^N \delta_{x^1}(dx^j)$, where μ is a compact-support Borel probability measure on F , and hence $\hat{\sigma}(\xi^1, \dots, \xi^N)$ is equal to $\hat{\mu}(\xi^1 + \cdots + \xi^N)$ for all $\xi^1, \dots, \xi^N \in \mathbf{R}^d$.

Therefore, Theorem 1.3 is a ready consequence of Corollaries 7.3 and 7.4. \square

8. Proof of Theorem 1.4

First, we elaborate on the connection between Theorem 1.4 and Bertoin's conjecture (Bertoin, 1999a, p. 49) that was mentioned briefly in the Introduction. Recall that a real-valued Lévy process is a *subordinator* if its sample functions are monotone a.s. (Bertoin, 1996; , 1999a; Fristedt, 1974; Sato, 1999).

Remark 8.1. We consider the special case that S_1 and S_2 are two (increasing) subordinators on \mathbf{R}_+ and $F := \{0\}$, and define two independent Lévy processes by $X_1 := S_1$ and $X_2 := -S_2$. Evidently, $X_1(t_1) + X_2(t_2) = 0$ for some $t_1, t_2 > 0$ if and only if $S_1(t_1) = S_2(t_2)$ for some $t_1, t_2 > 0$.

Let Σ_j denote the one-potential measure of S_j , and suppose $\Sigma_1(dx) = u_1(x) dx$, where u_1 is continuous on \mathbf{R} and strictly positive on $(0, \infty)$. Let U_j denote the one-potential measure of X_j . Then, $U_1 = \Sigma_1$ and $U_2 = \check{\Sigma}_2$, which we recall is the same as $\Sigma_2(-\bullet)$. It follows then that $U_1(dx) = u_1(x) dx$, and $(u_1 * U_2)(x) = \int_0^\infty u_1(x+y) \Sigma_2(dy)$ is strictly positive a.e. Therefore, we may apply Theorem 1.4 with $F := \{0\}$ to deduce that

$$(8.1) \quad \mathbf{P} \{S_1(t_1) = S_2(t_2) \text{ for some } t_1, t_2 > 0\} > 0 \quad \Leftrightarrow \quad Q(0) < \infty.$$

Because U_2 does not charge $(-\infty, 0)$ in the present setting,

$$(8.2) \quad Q(0) := \int_0^\infty \left[\frac{u_1(y) + u_1(-y)}{2} \right] U_2(dy) = \frac{1}{2} \int_0^\infty u_1(y) U_2(dy),$$

since $u_1(y) = 0$ for all $y < 0$. [In fact, Lemma 3.2 of Evans (1987b) tells us that $u_1(y) = 0$ for all $y \leq 0$.] Therefore, we conclude from (8.1) that

$$(8.3) \quad \mathbf{P} \{S_1(t_1) = S_2(t_2) \text{ for some } t_1, t_2 > 0\} = 0 \quad \Leftrightarrow \quad \int_0^\infty u_1(y) U_2(dy) = \infty.$$

Define the zero-potential measures U_j^0 as

$$(8.4) \quad U_j^0(A) := \mathbf{E} \left[\int_0^\infty \mathbf{1}_A(S_j(t)) dt \right],$$

for $j = 1, 2$ and Borel sets $A \subseteq \mathbf{R}_+$. Suppose $U_1^0(dx) = u^0(x) dx$, where u^0 is positive and continuous on $(0, \infty)$. Then, it is possible to adapt our methods, without any difficulties, to deduce also that

$$(8.5) \quad \mathbb{P} \{S_1(t_1) = S_2(t_2) \text{ for some } t_1, t_2 > 0\} = 0 \quad \Leftrightarrow \quad \int_0^\infty u_1^0(y) U_2^0(dy) = \infty.$$

The end of the proof of Theorem 1.5 contains a discussion which describes how similar changes can be made to adapt the proofs from statements about one-potentials U_j to those about zero-potentials U_j^0 . We omit the details, as they are not enlightening.

We mention the adaptations to zero-potentials of (8.5) for historical interest: (8.5) was conjectured by Bertoin, under precisely the stated conditions of this remark. Bertoin's conjecture was motivated in part by the fact (Bertoin, 1999b) that, under the very same conditions as above,

$$(8.6) \quad \mathbb{P} \{S_1(t_1) = S_2(t_2) \text{ for some } t_1, t_2 > 0\} = 0 \quad \Leftrightarrow \quad \sup_{z \in \mathbf{R}} \int_0^\infty u_1^0(y+z) U_2^0(dy) = \infty.$$

It is possible to deduce (8.6), and the same statement without the zero superscripts, from the present harmonic-analytic methods as well; see Lemma 5.4. The said extension goes well beyond the theory of subordinators, and is a sort of “low intensity maximum principle” (Salisbury, 1992). But we will not describe the details further, since we find the forms of (8.3) and (8.5) simpler to use, as well as easier to conceptualize. \square

Next we prove Theorem 1.4 without further ado.

Proof of Theorem 1.4. A direct computation reveals that the one-potential density of the two-parameter additive Lévy process $\mathfrak{X} := X_1 \oplus X_2$ is $u_1 * U_2$. Therefore, the equivalence of (1.14) and (1.15) follows from Part (c) in Proposition 6.6. Next we note that the one-potential density of \tilde{X}_1 is described by

$$(8.7) \quad v_1(x) := \frac{u_1(x) + u_1(-x)}{2} \quad \text{for all } x \in \mathbf{R}^d.$$

This function is positive a.e. and continuous away from the origin. We can use (6.15) and verify directly that $\tilde{\mathfrak{X}}$ has a one-potential density Q given by (1.17). Hence the last statement follows from Part (e) of Proposition 6.6. \square

9. Intersections of Lévy processes

The goal of this section is to prove Theorem 1.5. We first return briefly to Theorem 1.3 and discuss how it implies a necessary and sufficient condition for the existence of path-intersections for N independent Lévy processes. After proving Theorem 1.5, we conclude this section by presenting a nontrivial, though simple, example.

9.1. Existence of intersections. First we develop some general results on equilibrium measure that we believe might be of independent interest.

Choose and fix a compact set $F \subseteq \mathbf{R}^d$, and recall that $\text{cap}_{\Psi}(F)$ is the reciprocal of $\inf I_{\Psi}(\mu)$, where the infimum is taken over all Borel probability measures μ on F . It is not hard to see that when $\text{cap}_{\Psi}(F) > 0$, this infimum is in fact achieved. Indeed, for all $\epsilon > 0$ we can find a Borel probability measure μ_{ϵ} on F such that $I_{\Psi}(\mu_{\epsilon}) \leq (1 + \epsilon)/\text{cap}_{\Psi}(F)$. We can extract any subsequential weak limit μ of μ_{ϵ} 's. Evidently, μ is a Borel probability measure on F , and $I_{\Psi}(\mu) \leq 1/\text{cap}_{\Psi}(F)$ by Fatou's lemma. Because of the defining property of capacity, $I_{\Psi}(\mu) \geq 1/\text{cap}_{\Psi}(F)$ also. Therefore, $\text{cap}_{\Psi}(F)$ is in fact the reciprocal of the energy of μ .

Any Borel probability measure μ on F that has the preceding property is called an *equilibrium measure* on F . We now prove that there is only one equilibrium measure on a compact F .

Proposition 9.1. *If $F \subseteq \mathbf{R}^d$ is compact and has positive capacity $\text{cap}_{\Psi}(F)$, then there exists a unique Borel probability measure e_F on F such that $\text{cap}_{\Psi}(F) = 1/I_{\Psi}(e_F)$.*

Proof. For all finite signed probability measures μ and ν on \mathbf{R}^d define

$$(9.1) \quad I_{\Psi}(\mu, \nu) := \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \left(\frac{\overline{\hat{\mu}(\xi)} \hat{\nu}(\xi) + \overline{\hat{\nu}(\xi)} \hat{\mu}(\xi)}{2} \right) K_{\Psi}(\xi) d\xi.$$

This is well-defined, for example, if $I_{\Psi}(|\mu|) + I_{\Psi}(|\nu|) < \infty$, where $|\mu|$ is the total variation of μ . Indeed, we have the following Cauchy–Schwarz inequality: $|I_{\Psi}(\mu, \nu)|^2 \leq I_{\Psi}(|\mu|) \cdot I_{\Psi}(|\nu|)$. If σ is a finite [nonnegative] Borel measure on \mathbf{R}^d , then $I_{\Psi}(\sigma, \sigma)$ agrees with $I_{\Psi}(\sigma)$, and this is positive as long as σ is not the zero measure. However, we may note that slightly more general fact that if σ is a non-zero finite *signed* measure, then we still have

$$(9.2) \quad I_{\Psi}(\sigma, \sigma) > 0.$$

This follows from the fact that $K_{\Psi}(\xi) > 0$ for all $\xi \in \mathbf{R}^d$.

Let $c := 1/\text{cap}_{\Psi}(F)$, and suppose μ and ν were two distinct equilibrium measures on F . That is, $\mu \neq \nu$ but $I_{\Psi}(\mu) = I_{\Psi}(\nu) = 1/c$. In accord with (9.2),

$$(9.3) \quad 0 < I_{\Psi} \left(\frac{\mu - \nu}{2}, \frac{\mu - \nu}{2} \right) = \frac{c - I_{\Psi}(\mu, \nu)}{2}.$$

Consequently, $I_{\Psi}(\mu, \nu) < c$, and hence

$$\begin{aligned}
 (9.4) \quad I_{\Psi}\left(\frac{\mu + \nu}{2}\right) &= \frac{1}{4}I_{\Psi}(\mu) + \frac{1}{4}I_{\Psi}(\nu) + \frac{1}{2}I_{\Psi}(\mu, \nu) \\
 &= \frac{c + I_{\Psi}(\mu, \nu)}{2} \\
 &< c.
 \end{aligned}$$

Because $\frac{1}{2}(\mu + \nu)$ is a Borel probability measure on F , this contradicts the fact that c is the smallest possible energy on F . \square

Next we note with the following computation of the equilibrium measure in a specific class of examples. The following result is related quite closely to the celebrated local ergodic theorem of Csiszár (1965). [We hope to elaborate on this connection elsewhere.] See also Proposition A3 of Khoshnevisan, Xiao, and Zhong (2003) for a related result.

Proposition 9.2. *Suppose F is a fixed compact subset of \mathbf{R}^d with a nonvoid interior. If $\text{cap}_{\Psi}(F) > 0$, then e_F is the normalized Lebesgue measure on F .*

Proof. Our strategy is to prove that e_F is translation invariant.

Let $\mathcal{R}(r)$ denote the collection of all closed “upright” cubes of the form

$$(9.5) \quad I := [s_1, s_1 + r] \times \cdots \times [s_d, s_d + r] \subseteq F,$$

such that all of the s_i ’s are rational numbers and $r > 0$. Each $\mathcal{R}(r)$ is a countable collection, and hence we can (and will) enumerate its elements as $I_1(r), I_2(r), \dots$.

For all $i, j \geq 1$ and $r > 0$ we choose and fix a one-to-one onto piecewise-linear map $\theta_{i,j,r} : F \rightarrow F$ that has the following properties:

- If $a \notin I_i(r) \cup J_j(r)$, then $\theta_{i,j,r}(a) = a$;
- $\theta_{i,j,r}$ maps $I_i(r)$ onto $J_j(r)$ bijectively; and
- $\theta_{i,j,r}$ maps $J_j(r)$ onto $I_i(r)$ bijectively.

To be concrete, let us write $I_i(r) = I_j(r) + b$, where $b \in \mathbf{R}^d$. Then we define

$$(9.6) \quad \theta_{i,j,r}(a) = \begin{cases} a - b & \text{if } a \in I_i(r), \\ a + b & \text{if } a \in J_j(r), \\ a & \text{if } a \notin I_i(r) \cup J_j(r). \end{cases}$$

It can be verified that $\theta_{i,j,r} \circ \theta_{i,j,r} = id$, the identity map.

For all integers $n \geq 1$ and $r > 0$ consider

$$(9.7) \quad \rho_{n,r} := \frac{1}{n^2} \sum_{1 \leq i, j \leq n} (e_F \circ \theta_{i,j,r}^{-1}).$$

Obviously, each $\rho_{n,r}$ is a probability measure on F , and

$$(9.8) \quad \rho_{n,r}(I_i(r)) = \rho_{n,r}(I_j(r)) \quad \text{for all } 1 \leq i, j \leq n.$$

Furthermore, a direct computation reveals that for all $i, j \geq 1$ and $r > 0$,

$$(9.9) \quad \left| \widehat{e_F \circ \theta_{i,j,r}^{-1}} \right| = |\widehat{e_F}| \quad \text{pointwise.}$$

Therefore, we can apply Minkowski's inequality, to the norm $\mu \mapsto \sqrt{I_\Psi(\mu)}$, to find that

$$(9.10) \quad \begin{aligned} \sqrt{I_\Psi(\rho_{n,r})} &\leq \frac{1}{n^2} \sum_{1 \leq i, j \leq n} \sqrt{I_\Psi(e_F \circ \theta_{i,j,r}^{-1})} \\ &= \sqrt{I_\Psi(e_F)}. \end{aligned}$$

By the uniqueness of equilibrium measure (Proposition 9.1), $e_F = \rho_{n,r}$ for all $n \geq 1$ and rationals $r > 0$. This and (9.8) together prove that $e_F(I) = e_F(J)$ for all $r > 0$ and all $I, J \in \mathcal{R}(r)$. A monotone-class argument reveals that for all Borel sets $I \subset F$ and $b \in \mathbf{R}^d$, $e_F(I) = e_F(b+I)$, provided that $I, b+I \subseteq F$. Because the Lebesgue measure is characterized by its translation invariance, this implies the proposition. \square

Lemma 5.6 and Proposition 9.2, and Theorem 1.3 together imply the following variant of a theorem of Fitzsimmons and Salisbury (1989).

Corollary 9.3. *Let X_1, \dots, X_N be independent Lévy processes on \mathbf{R}^d , and assume that each X_j has a one-potential density u_j that is continuous on \mathbf{R}^d , positive at zero, and finite on $\mathbf{R}^d \setminus \{0\}$. Then,*

$$(9.11) \quad \mathbf{P} \{X_1(t_1) = \dots = X_N(t_N) \text{ for some } t_1, \dots, t_N > 0\} > 0$$

if and only if

$$(9.12) \quad \prod_{j=1}^N \left(\frac{u_j(\bullet) + u_j(-\bullet)}{2} \right) \in L_{loc}^1(\mathbf{R}^d).$$

Proof. Clearly, (9.11) holds if and only if there exists $n > 0$ such that

$$(9.13) \quad \mathbf{P} \{X_1(t_1) = \dots = X_N(t_N) \in [-n, n]^d \text{ for some } t_1, \dots, t_N > 0\} > 0.$$

Theorem 1.3 implies that (9.13) holds if and only if $\text{cap}_\Psi([-n, n]^d) > 0$. Lemma 5.6 and Proposition 9.2 together prove the result. \square

9.2. Proof of Theorem 1.5. We begin by proving the first part; thus, we assume only that the u_j 's exist and are a.e.-positive.

Let \mathfrak{S} denote an independent M -parameter additive stable process of index $\alpha \in (0, 2)$; see (10.9). Next, we consider the $(N + M)$ -parameter process $\mathfrak{Y} := \otimes_{j=1}^N (X_j - \mathfrak{S})$; i.e.,

$$(9.14) \quad \mathfrak{Y}(\mathbf{s} \otimes \mathbf{t}) := \begin{pmatrix} X_1(s_1) - \mathfrak{S}(\mathbf{t}) \\ \vdots \\ X_N(s_N) - \mathfrak{S}(\mathbf{t}) \end{pmatrix} \quad \text{for all } \mathbf{s} \in \mathbf{R}_+^N, \mathbf{t} \in \mathbf{R}_+^M.$$

It is not hard to adapt the discussion of the first few paragraphs in §7 to the present situation and deduce that \mathfrak{Y} is an $(N + M)$ -parameter additive Lévy process, with values in $(\mathbf{R}^d)^N$, and that for all $\mathbf{s} \in \mathbf{R}_+^N$, $\mathbf{t} \in \mathbf{R}_+^M$, and $\xi := \xi^1 \otimes \cdots \otimes \xi^N \in (\mathbf{R}^d)^N$,

$$(9.15) \quad \mathbb{E} \exp(i\xi \cdot \mathfrak{Y}(\mathbf{s} \otimes \mathbf{t})) = \exp \left(- \sum_{k=1}^N s_k \Psi_k(\xi^k) - \sum_{l=1}^M t_l \|\xi^1 + \cdots + \xi^N\|^\alpha \right).$$

We can conclude readily from this that the characteristic exponent of \mathfrak{Y} is defined by

$$(9.16) \quad \Theta(\xi) := \left(\Psi_1(\xi^1), \dots, \Psi_N(\xi^N), \underbrace{\left\| \sum_{j=1}^N \xi^j \right\|^\alpha, \dots, \left\| \sum_{j=1}^N \xi^j \right\|^\alpha}_{M \text{ times}} \right),$$

for all $\xi := \xi^1 \otimes \cdots \otimes \xi^N \in (\mathbf{R}^d)^N$.

It follows readily from this and Lemma 7.2 that \mathfrak{Y} and $\tilde{\mathfrak{Y}}$ both have positive one-potential densities. Moreover, a direct computation involving the inversion formula reveals that the potential density of $\tilde{\mathfrak{Y}}$ is defined by

$$(9.17) \quad v(x^1 \otimes \cdots \otimes x^N) = \int_{\mathbf{R}^d} \prod_{j=1}^N \left(\frac{u_j(x^j - y) + u_j(y - x^j)}{2} \right) w(y) dy,$$

for all $x := (x^1 \otimes \cdots \otimes x^N) \in (\mathbf{R}^d)^N$. Here, w denotes the one-potential density of \mathfrak{S} . That is,

$$(9.18) \quad w(y) := \int_{\mathbf{R}_+^M} p_{\mathbf{t}}(y) e^{-[\mathbf{t}]} d\mathbf{t},$$

where $p_{\mathbf{t}}$ denotes the density of $\mathfrak{S}(\mathbf{t})$. That is, $p_{\mathbf{t}}(y) := (2\pi)^{-d} \int_{\mathbf{R}^d} \exp(-iy \cdot z - [\mathbf{t}]\|z\|^\alpha) dz$ for all $\mathbf{t} \in \mathbf{R}_+^M$ and $y \in \mathbf{R}^d$.

Since the u_j 's are everywhere positive (Lemma 7.2), Proposition 6.6 tell us that $0 \in \mathfrak{Y}(\mathbf{R}_+^N \times \mathbf{R}_+^M)$ with positive probability if and only if $\text{cap}_{\Theta}(\{0\}) > 0$. Thus, the preceding

positive capacity condition is equivalent to the integrability of the function K_{Θ} . That is,

$$(9.19) \quad \begin{aligned} & 0 \in \mathfrak{Y}(\mathbf{R}_+^N \times \mathbf{R}_+^M) \quad \text{with positive probability} \\ \Leftrightarrow & \int_{(\mathbf{R}^d)^N} \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi^j)} \right) \frac{d\xi}{1 + \|\xi^1 + \dots + \xi^N\|_{\alpha M}} < \infty. \end{aligned}$$

On the other hand, it is manifestly the case that $\mathfrak{Y}(\mathbf{R}_+^N \times \mathbf{R}_+^M)$ contains the origin if and only if the intersection of $\cap_{j=1}^N X_j(\mathbf{R}_+)$ and $\mathfrak{S}(\mathbf{R}_+^M)$ is nonempty. Thanks to Theorem 4.4.1 of Khoshnevisan (2002, p. 428), for all Borel sets $F \subseteq \mathbf{R}^d$, $\mathbf{P}\{\mathfrak{S}(\mathbf{R}_+^M) \cap F \neq \emptyset\} > 0$ if and only if $\mathcal{C}_{d-M\alpha}(F) > 0$, provided that we also assume that $d > M\alpha$. Suppose, then, that $d > M\alpha$. We can apply the preceding, conditionally on X_1, \dots, X_N , and deduce that

$$(9.20) \quad \begin{aligned} & 0 \in \mathfrak{Y}(\mathbf{R}_+^N \times \mathbf{R}_+^M) \quad \text{with positive probability} \\ \Leftrightarrow & \mathcal{C}_{d-M\alpha} \left(\bigcap_{j=1}^N X_j(\mathbf{R}_+) \right) > 0 \quad \text{with positive probability.} \end{aligned}$$

We compare the preceding display to (9.19), and choose M and $\alpha \in (0, (d/M) \wedge 2)$, such that $d - M\alpha$ is any prescribed number $s \in (0, d)$. This yields the following: For all $s \in (0, d)$,

$$(9.21) \quad \begin{aligned} & \mathcal{C}_s \left(\bigcap_{j=1}^N X_j(\mathbf{R}_+) \right) > 0 \quad \text{with positive probability} \\ \Leftrightarrow & \int_{(\mathbf{R}^d)^N} \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi^j)} \right) \frac{d\xi}{1 + \|\xi^1 + \dots + \xi^N\|^{d-s}} < \infty. \end{aligned}$$

It follows from (9.21), Frostman's theorem (Khoshnevisan, 2002, p. 521) and an argument similar to the proof of Theorem 3.2 in Khoshnevisan, Shieh and Xiao (2007) that the first identity of (1.18) holds almost surely on $\{\cap_{j=1}^N X_j(\mathbf{R}_+) \neq \emptyset\}$.

In order to obtain (1.19), we assume also that the u_j 's are continuous on \mathbf{R}^d and finite on $\mathbf{R}^d \setminus \{0\}$. Thanks to Proposition 6.6, $\mathfrak{Y}(\mathbf{R}_+^N \times \mathbf{R}_+^M)$ contains zero with positive probability if and only if $v(0) < \infty$. Thus, (9.17) and (9.20) imply that

$$(9.22) \quad \begin{aligned} & \mathcal{C}_{d-M\alpha} \left(\bigcap_{j=1}^N X_j(\mathbf{R}_+) \right) > 0 \quad \text{with positive probability} \\ \Leftrightarrow & \int_{\mathbf{R}^d} \prod_{j=1}^N \left(\frac{u_j(y) + u_j(-y)}{2} \right) w(y) dy < \infty. \end{aligned}$$

If we could replace $w(y)$ by $\|y\|^{-d+\alpha M}$, then we could finish the proof by choosing M and α suitably, and then appealing to the Frostman theorem and the argument in Khoshnevisan,

Shieh and Xiao (2007). [This is how we completed the proof of the first part of the proof as well.] Thus, our goal is to derive (9.22).

Unfortunately, it is not possible to simply replace $w(y)$ by $\|y\|^{-d+\alpha M}$ by simple real-variable arguments. Nonetheless, we recall the following fact from Khoshnevisan (2002, Exercise 4.1.4, p. 423): There exist $c, C \in (0, \infty)$ such that $c\|y\|^{-d+M\alpha} \leq w(y) \leq C\|y\|^{-d+M\alpha}$ for all $y \in (-1, 1)^d$. Moreover, the upper bound holds for all $y \in \mathbf{R}^d$ (*loc. cit.*, Eq. (2), p. 423); the lower bound [provably] does not. It follows then that

$$\begin{aligned}
 (9.23) \quad & \int_{\mathbf{R}^d} \prod_{j=1}^N \left(\frac{u_j(y) + u_j(-y)}{2} \right) \frac{dy}{\|y\|^{d-M\alpha}} < \infty \\
 & \Rightarrow \mathcal{C}_{d-M\alpha} \left(\bigcap_{j=1}^N X_j(\mathbf{R}_+) \right) > 0 \quad \text{with positive probability} \\
 & \Rightarrow \int_{(-1,1)^d} \prod_{j=1}^N \left(\frac{u_j(y) + u_j(-y)}{2} \right) \frac{dy}{\|y\|^{d-M\alpha}} < \infty.
 \end{aligned}$$

If we could replace $(-1, 1)^d$ by \mathbf{R}^d in the last display, then our proof follows the outline mentioned earlier. Thus, we merely point out how to derive (9.22), with $w(y)$ replaced by $\|y\|^{-d+M\alpha}$, and omit the remainder of the argument.

The proof hinges on a few modifications to the entire theory outlined here. We describe them [very] briefly, since it is easy—though tedious—to check that the present changes go through unhindered.

Define the operator $R^{1\otimes 0}$ via

$$(9.24) \quad (R^{1\otimes 0}f)(x) := \frac{1}{2^N} \mathbb{E} \left[\int_{\mathbf{R}^N \times \mathbf{R}^M} f(x + \mathfrak{Y}(\mathbf{s} \otimes \mathbf{t})) e^{-[\mathbf{s}]} d\mathbf{s} d\mathbf{t} \right] \quad \text{for all } x \in (\mathbf{R}^d)^N.$$

If, we replace $\exp(-[\mathbf{s}])$ by $\exp(-[\mathbf{s}] - [\mathbf{t}])$, then $R^{1\otimes 0}f$ turns into Rf .

As is, the operator $R^{1\otimes 0}f$ fails to map $L^p((\mathbf{R}^d)^N)$ into $L^p((\mathbf{R}^d)^N)$. But this is a minor technical nuisance, since one can check directly that

$$(9.25) \quad (R^{1\otimes 0}f)(x) = \int_{\mathbf{R}^d} f(x+y) v^{1\otimes 0}(y) dy,$$

where $v^{1\otimes 0}$ is defined exactly as v was, but with $w(y)$ replaced by a certain constant times $\|y\|^{-d+M\alpha}$. And this shows fairly readily that if f is a compactly-supported function in $L^1((\mathbf{R}^d)^N)$, then $R^{1\otimes 0}f \in L^1_{loc}((\mathbf{R}^d)^N)$. To complete our proof, we need to redevelop the potential theory of the additive Lévy process \mathfrak{Y} , but this time in terms of $R^{1\otimes 0}f$ and $v^{1\otimes 0}$.

The fact that $R^{1\otimes 0} : L^1_c((\mathbf{R}^d)^N) \rightarrow L^1_{loc}((\mathbf{R}^d)^N)$ provides sufficient regularity to allow us to push the Fourier analysis of the present paper through without change. And the end result

is that $\mathfrak{Y}(\mathbf{R}_+^N \times \mathbf{R}_+^M)$ contains zero if and only if $v^{1 \otimes 0}(0)$ is finite. Now we can complete the proof, but with $v^{1 \otimes 0}$ in place of v everywhere. \square

9.3. An example. Suppose X_1, \dots, X_N are independent isotropic stable processes in \mathbf{R}^d with respective Fourier transforms $E \exp(i\xi \cdot X_j(t)) = \exp(-c_j t \|\xi\|^{\alpha_j})$ for all $t \geq 0$, $\xi \in \mathbf{R}^d$, and $1 \leq j \leq N$. Here, c_1, \dots, c_N are constants, and $0 < \alpha_1, \dots, \alpha_N < 2 \wedge d$ are the indices of stability. We are primarily interested in the case that $N \geq 2$, but the following remarks apply to the case $N = 1$ equally well.

In this section we work out some of intersection properties of X_1, \dots, X_N . It is possible to construct much more sophisticated examples. We study the present setting because it provides us with the simplest nontrivial example of its type.

It is known that each X_j has a continuous positive one-potential density u_j , and there exist $c, C \in (0, \infty)$ such that $c\|x\|^{-d+\alpha_j} \leq u_j(x) \leq C\|x\|^{-d+\alpha_j}$ for all $x \in (-1, 1)^d$ and $1 \leq j \leq N$. [These assertions follow, for example, from Corollary 3.2.1 on page 379, and Lemma 3.4.1 on page 383 of Khoshnevisan (2002).] Consequently, Theorem 1.3 immediately implies that for all Borel sets $F \subseteq \mathbf{R}^d$,

$$(9.26) \quad \mathbb{P} \left\{ \bigcap_{j=1}^N X_j(\mathbf{R}_+) \cap F \neq \emptyset \right\} > 0 \quad \Leftrightarrow \quad \mathcal{C}_{Nd - \sum_{j=1}^N \alpha_j}(F) > 0.$$

In the case that $N = 2$, a slightly more general form of this was found in Khoshnevisan (2002, Theorem 4.4.1, p. 428) by using other methods. We may apply the preceding with $F = \mathbf{R}^d$, and appeal to Taylor's theorem (*loc. cit.*, Corollary 2.3.1, p. 525) to find that

$$(9.27) \quad \mathbb{P} \left\{ \bigcap_{k=1}^N X_k(\mathbf{R}_+) \neq \emptyset \right\} > 0 \quad \Leftrightarrow \quad (N-1)d < \sum_{j=1}^N \alpha_j.$$

Theorem 1.5, and a direct computation in polar coordinates, together show that the slack in the preceding inequality determines the Hausdorff dimension of the set $\bigcap_{k=1}^N X_k(\mathbf{R}_+)$ of intersection points. That is, almost surely on $\left\{ \bigcap_{k=1}^N X_k(\mathbf{R}_+) \neq \emptyset \right\}$,

$$(9.28) \quad \dim_{\mathbf{H}} \bigcap_{k=1}^N X_k(\mathbf{R}_+) = \left[\sum_{k=1}^N \alpha_k - (N-1)d \right]_+.$$

This formula continues to hold in case some, or even all, of the α_j 's are equal to or exceed d . We omit the details.

9.4. Remarks on multiple points. Next we mention how the preceding fits in together with the well-known conjecture of Hendricks and Taylor (1979) that was solved in Fitzsimmons and Salisbury (1989), and also make a few related remarks.

Remark 9.4. Recall that a [single] Lévy process X with values in \mathbf{R}^d has N -multiple points if and only if there exist times $0 < t_1 < \dots < t_N < \infty$ such that $X(t_1) = X(t_2) = \dots = X(t_N)$. [We are ruling out the possibility that $t_1 = 0$ merely to avoid degeneracies.]

By localization and the Markov property, X has N -multiple points almost surely if and only if there are times $t_1, \dots, t_N \in (0, \infty)$ such that $X_1(t_1) = \dots = X_N(t_N)$ with positive probability, where X_1, \dots, X_N are N i.i.d. copies of X .

Suppose that X has a continuous and positive [equivalently, positive-at-zero] one-potential density u that is finite on $\mathbf{R}^d \setminus \{0\}$. Then according to Corollary 9.3, X has N -multiple points if and only if $u(\bullet) + u(-\bullet) \in L_{loc}^N(\mathbf{R}^d)$. An application of Hölder's inequality reveals then that X has N -multiple points if and only if $u \in L_{loc}^N(\mathbf{R}^d)$. This is more or less the well-known condition of Hendricks and Taylor (1979).

More generally, the following can be deduced with no extra effort: Under the preceding conditions, given a nonrandom Borel set $F \subseteq \mathbf{R}^d$,

$$(9.29) \quad \mathbf{P} \{ \text{there exist } 0 < t_1 < \dots < t_N \text{ such that } X(t_1) = \dots = X(t_N) \in F \} > 0$$

if and only if there exists a compact-support Borel probability measure μ on F such that

$$(9.30) \quad \iint [u(x - y)]^N \mu(dx) \mu(dy) < \infty.$$

See, for example, Theorem 5.1 of Fitzsimmons and Salisbury (1989). \square

Remark 9.5. We mention the following formula for the Hausdorff dimension of the N -multiple points of a Lévy process X : Under the conditions stated in Remark 9.4,

$$(9.31) \quad \dim_{\mathbf{H}} M_N = \sup \left\{ s \in (0, d) : \int_{\mathbf{R}^d} \frac{[u(z)]^N}{\|z\|^s} dz < \infty \right\} \quad \text{a.s.},$$

where M_N denotes the collection of all N -multiple points. That is, M_N is the set of all $x \in \mathbf{R}^d$ for which the cardinality of $X^{-1}(\{x\})$ is at least N . This formula appears to be new. Hawkes (1978b, Theorem 2) contains a similar formula—with $\int_{\mathbf{R}^d}$ replaced by $\int_{(-1,1)^d}$ —which is shown to be valid for all isotropic [spherically symmetric] Lévy processes that have measurable transition densities.

In order to prove (9.31), we appeal to Theorem 1.5 and the Markov property to first demonstrate that the dimension formula in (9.31) is valid almost surely on $\{M_N \neq \emptyset\}$. Then the conclusion follows from the fact that the event $\{M_N \neq \emptyset\}$ satisfies a zero-one law; that zero-one law is itself proved by adapting the argument of Orey (1967, p. 124) or Evans (1987b, pp. 365–366). We omit the details as the method is nowadays considered standard. \square

Remark 9.6. This is a natural place to complete a computation that we alluded to earlier in Open Problem 4. Namely, we wish to prove that under the preceding conditions on the Lévy process X ,

$$(9.32) \quad \dim_{\mathbb{H}} M_2 = \sup \left\{ s \in (0, d) : \int_{(-1,1)^d} \frac{[u(z)]^2}{\|z\|^s} dz < \infty \right\} \quad \text{a.s.},$$

By the Markov property, it suffices to derive (9.32) with M_2 replaced by $X_1(\mathbf{R}_+) \cap X_2(\mathbf{R}_+)$, where X_1 and X_2 are independent copies of X .

We have seen already that if $u \notin L^2_{loc}(\mathbf{R}^d)$, then $X_1(\mathbf{R}_+) \cap X_2(\mathbf{R}_+) = \emptyset$, and there is nothing left to prove. Thus, we may consider only the case that $u \in L^2_{loc}(\mathbf{R}^d)$. In this case, $X_1(\mathbf{R}_+) \cap X_2(\mathbf{R}_+) \neq \emptyset$, which as we have seen is equivalent to the condition that $(X_1 \ominus X_2)(\mathbf{R}^2_+)$ contains the origin. This and Theorem 1.1 together prove that $K_{\Psi} \in L^2(\mathbf{R}^d)$, where $K_{\Psi}(\xi) = \text{Re}(1 + \Psi(\xi))^{-1}$ and $\mathbb{E} \exp(i\xi \cdot X(t)) = \exp(-t\Psi(\xi))$ in the present case. Now Lemma 6.2 shows that \hat{v} —and hence v —is square integrable, where $v := \frac{1}{2}(u(\bullet) + u(-\bullet))$. That is, $u \in L^2(\mathbf{R}^d)$. From here, it is a simple matter to check that (9.31) with $N = 2$ implies (9.32). \square

10. Zero-one laws

We conclude this paper by deriving two zero-one laws: One for the Lebesgue measure of the range of additive Lévy processes; and another for capacities of the range of additive Lévy processes.

The following was proved first in Khoshnevisan, Xiao, and Zhong (2003) under a mild technical condition. Here we remove the technical condition (1.3) of that paper, and derive this result as an elementary consequence of Theorem 1.1.

Proposition 10.1. *Let \mathfrak{X} be a general N -parameter additive Lévy process in \mathbf{R}^d with exponent Ψ . Then,*

$$(10.1) \quad \mathbb{E} [\lambda_d(\mathfrak{X}(\mathbf{R}^N_+))] > 0 \quad \text{if and only if} \quad \int_{\mathbf{R}^d} \prod_{j=1}^N \text{Re} \left(\frac{1}{1 + \Psi_j(\xi)} \right) d\xi < \infty.$$

Suppose, in addition, that $\tilde{\mathfrak{X}}$ has a one-potential density that is continuous away from the origin, and \mathfrak{X} has an a.e.-positive one-potential density. Then, with probability one, $\lambda_d(\mathfrak{X}(\mathbf{R}^N_+))$ is zero or infinity.

Proof. According to Theorem 1.1, $\mathbb{E}[\lambda_d(\mathfrak{X}(\mathbf{R}^N_+))]$ is positive if and only if $\text{cap}_{\Psi}(\{0\}) > 0$. But the only probability measure on $\{0\}$ is δ_0 , and hence $\text{cap}_{\Psi}(\{0\}) > 0$ iff $K_{\Psi} \in L^1(\mathbf{R}^d)$. This is the integrability condition of (10.1), and implies the assertion made in (10.1).

Now the proof of Proposition 6.5 of Khoshnevisan et al. (2003) goes through unhindered to conclude the remainder of the proof. We include it here for the sake of completeness. Throughout the rest of this proof, we assume that \mathfrak{X} has an a.e.-positive one-potential density.

Our task is to prove that if $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))]$ is finite, then it is zero. Note that for all $n > 0$,

$$(10.2) \quad \begin{aligned} E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))] \\ \geq E[\lambda_d(\mathfrak{X}([0, n]^N))] + E[\lambda_d(\mathfrak{X}([n, \infty)^N))] - E[\lambda_d(\mathfrak{X}([0, n]^N) \cap \mathfrak{X}'(\mathbf{R}_+^N))], \end{aligned}$$

where \mathfrak{X}' is an independent copy of \mathfrak{X} . Therefore, if $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))]$ is finite, then

$$(10.3) \quad E[\lambda_d(\mathfrak{X}([0, n]^N))] \leq E[\lambda_d(\mathfrak{X}([0, n]^N) \cap \mathfrak{X}'(\mathbf{R}_+^N))].$$

The inequality is, in fact, an equality. Let $n \uparrow \infty$ to deduce that if $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))]$ is finite, then $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))] = E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N) \cap \mathfrak{X}'(\mathbf{R}_+^N))]$. Consider the function

$$(10.4) \quad \phi(a) := P\{a \in \mathfrak{X}(\mathbf{R}_+^N)\} \quad \text{for all } a \in \mathbf{R}^d.$$

Then, we have just proved that $\int_{\mathbf{R}^d} \phi(a) da = \int_{\mathbf{R}^d} |\phi(a)|^2 da$. Since $0 \leq \phi(a)(1 - \phi(a)) \leq 1$ for all $a \in \mathbf{R}^d$, it follows that $\phi \in \{0, 1\}$ almost everywhere. Consequently, if $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))]$ is finite, then

$$(10.5) \quad E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))] = \lambda_d(\phi^{-1}(\{1\})).$$

According to Proposition 6.6, either $\phi(a) = 0$ for all a , or $\phi(a) = 1$ for all a . Since $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))] < \infty$, this and (10.5) together prove that $\phi(a) = 0$ for all a , and hence $E[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N))] = 0$. \square

For all $s > 0$ define $\mathcal{C}_s(A)$ to be the s -dimensional *Bessel–Riesz capacity* of the Borel set $A \subseteq \mathbf{R}^d$. That is,

$$(10.6) \quad \mathcal{C}_s(A) := \left[\inf_{\mu \in \mathcal{P}_c(A)} \iint \frac{\mu(dx) \mu(dy)}{\|x - y\|^s} \right]^{-1},$$

where $\inf \emptyset := \infty$, $1/\infty := 0$, and we recall that $\mathcal{P}_c(A)$ denotes the collection of all compact-support Borel probability measures on A . Thus, $\mathcal{C}_s = \mathcal{C}_{\kappa_{d-s}}$, where $\kappa_\alpha(x) := \|x\|^{-d+\alpha}$ denotes the $(d - \alpha)$ -dimensional Riesz kernel. The goal of this subsection is to derive a zero-one law for the capacity of the range of an arbitrary additive Lévy process \mathfrak{X} .

Proposition 10.2. *If \mathfrak{X} denotes an N -parameter additive Lévy process in \mathbf{R}^d , then for all $\beta \in (0, d)$ fixed, the chances are either zero or one that $\mathcal{C}_\beta(\mathfrak{X}(\mathbf{R}_+^N))$ is strictly positive.*

The following formula for the Hausdorff dimension of $\mathfrak{X}(\mathbf{R}_+^N)$ is an immediate application of Proposition 10.2 and the methods of Khoshnevisan and Xiao (2004, Theorem 4.1): If \mathfrak{X} is

an additive Lévy process with values in \mathbf{R}^d with characteristic exponent (Ψ_1, \dots, Ψ_N) , then almost surely,

$$(10.7) \quad \dim_{\mathbf{H}}(\mathfrak{X}(\mathbf{R}_+^N)) = \sup \left\{ \beta \in (0, d) : \int_{\mathbf{R}^d} \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi)} \right) \frac{d\xi}{\|\xi\|^{d-\beta}} < +\infty \right\}.$$

Here $\sup \emptyset := 0$.

Proof of Proposition 10.2. We choose and fix a $\beta \in (0, d)$, and assume that

$$(10.8) \quad \mathbf{P}\{\mathcal{C}_\beta(\mathfrak{X}(\mathbf{R}_+^N)) > 0\} > 0,$$

for there is nothing to prove otherwise.

Let us choose an integer $M \geq 1$ and a real number $\alpha \in (0, 2]$ such that $\beta = d - M\alpha$. After enlarging the probability space, if need be, we may introduce M i.i.d. isotropic stable processes S_1, \dots, S_M —independent also of all X_j 's—such that each S_j has stability index α . In this way, we can consider also the M -parameter additive Lévy process

$$(10.9) \quad \mathfrak{S}(\mathbf{u}) := S_1(u_1) + \dots + S_M(u_M) \quad \text{for all } \mathbf{u} := (u_1, \dots, u_M) \in \mathbf{R}_+^M.$$

defined on \mathbf{R}^d .

According to Theorem 7.2 of Khoshnevisan, Xiao, and Zhong (2003),

$$(10.10) \quad \mathbf{P}\{\mathcal{C}_\beta(\mathfrak{X}(\mathbf{R}_+^N)) > 0\} > 0 \Leftrightarrow \mathbf{E}[\lambda_d(\mathfrak{X}(\mathbf{R}_+^N) \oplus \mathfrak{S}(\mathbf{R}_+^M))] > 0.$$

Because $\mathfrak{X} \oplus \mathfrak{S}$ is itself an $(N+M)$ -parameter additive Lévy process, Proposition 10.1 implies that the right-hand side of (10.10) is equivalent to

$$(10.11) \quad \int_{\mathbf{R}^d} \left(\frac{1}{1 + \|\xi\|^\alpha} \right)^M \prod_{j=1}^N \operatorname{Re} \left(\frac{1}{1 + \Psi_j(\xi)} \right) d\xi < \infty.$$

It remains to prove $\mathbf{P}\{\mathcal{C}_\beta(\mathfrak{X}(\mathbf{R}_+^N)) > 0\} = 1$.

The proof is similar to that of Proposition 2.8 in Khoshnevisan and Xiao (2005). Define the random probability measure m on $\mathfrak{X}(\mathbf{R}_+^N)$ by

$$(10.12) \quad \int_{\mathbf{R}^d} f(x) m(dx) = \int_{\mathbf{R}_+^N} f(\mathfrak{X}(\mathbf{t})) e^{-[\mathbf{t}]} d\mathbf{t},$$

where $f : \mathbf{R}^d \rightarrow \mathbf{R}_+$ denotes a Borel measurable function. It follows from Lemma 5.6 and the Fubini-Tonelli theorem that

$$(10.13) \quad \mathbf{E} \left[\iint \frac{m(dx) m(dy)}{\|x - y\|^\beta} \right] = \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} \mathbf{E}(\|\hat{m}(\xi)\|^2) \frac{d\xi}{\|\xi\|^{d-\beta}}.$$

We may observe that

$$\begin{aligned}
 \mathbb{E} (\|\hat{m}(\xi)\|^2) &= \int_{\mathbf{R}_+^N} \int_{\mathbf{R}_+^N} \mathbb{E} [e^{-i\xi \cdot (\mathfrak{X}(\mathbf{t}) - \mathfrak{X}(\mathbf{s}))}] e^{-[\mathbf{s}] - [\mathbf{t}]} d\mathbf{s} d\mathbf{t} \\
 (10.14) \qquad &= \prod_{j=1}^N \left[\int_0^\infty \int_0^\infty e^{-s_j - t_j - |s_j - t_j| \Psi_j(\text{sgn}(s_j - t_j)\xi)} ds_j dt_j \right].
 \end{aligned}$$

A direct computation shows that the preceding is equal to $K_{\Psi}(\xi)$; see also the proof of Lemma 3.4. Thanks to this and (10.11), the final integral in (10.13) is finite, whence it follows that $\mathcal{C}_{\beta}(\mathfrak{X}(\mathbf{R}_+^N)) > 0$ almost surely. \square

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